

Dynamic Large Deviations: In search of Matching Bounds

Kac's Process and the Zero-Range Process

D. Heydecker, partially joint work with Benjamin Gess



Slides: [danielheydecker.wordpress.com](#) → Research

Max-Planck Institut für Mathematik in den Naturwissenschaften, Leipzig

Introduction

Introduction and Problem

- Starting point:

Introduction and Problem

- Starting point: Kac's model of the spatially homogeneous Boltzmann equation $\mu_t^N \rightarrow \mu_t$

Introduction and Problem

- Starting point: Kac's model of the spatially homogeneous Boltzmann equation $\mu_t^N \rightarrow \mu_t$, Zero-range process $\eta_t^N \rightarrow u_t$:

Introduction and Problem

- Starting point: Kac's model of the spatially homogeneous Boltzmann equation $\mu_t^N \rightarrow \mu_t$, Zero-range process $\eta_t^N \rightarrow u_t$: microscopic models for the nonlinear PDEs

Introduction and Problem

- Starting point: Kac's model of the spatially homogeneous Boltzmann equation $\mu_t^N \rightarrow \mu_t$, Zero-range process $\eta_t^N \rightarrow u_t$: microscopic models for the nonlinear PDEs

$$\partial_t \mu_t = Q(\mu_t, \mu_t); \quad \partial_t u_t = \frac{1}{2} \Delta \Phi(u_t).$$

Introduction and Problem

- Starting point: Kac's model of the spatially homogeneous Boltzmann equation $\mu_t^N \rightarrow \mu_t$, Zero-range process $\eta_t^N \rightarrow u_t$: microscopic models for the nonlinear PDEs

$$\partial_t \mu_t = Q(\mu_t, \mu_t); \quad \partial_t u_t = \frac{1}{2} \Delta \Phi(u_t).$$

- Seek to characterise the large deviations around these limits:

$$\mathbb{P}(\mu_\bullet^N \approx \mu_\bullet) \sim \exp(-N \mathcal{I}_{\text{KP}}(\mu_\bullet));$$

$$\mathbb{P}(\eta_\bullet^N \approx u_\bullet) \sim \exp(-N^d \mathcal{I}_{\text{ZRP}}(u_\bullet))$$

Introduction and Problem

- Starting point: Kac's model of the spatially homogeneous Boltzmann equation $\mu_t^N \rightarrow \mu_t$, Zero-range process $\eta_t^N \rightarrow u_t$: microscopic models for the nonlinear PDEs

$$\partial_t \mu_t = Q(\mu_t, \mu_t); \quad \partial_t u_t = \frac{1}{2} \Delta \Phi(u_t).$$

- Seek to characterise the large deviations around these limits:

$$\mathbb{P}(\mu_{\bullet}^N \approx \mu_{\bullet}) \sim \exp(-N \mathcal{I}_{\text{KP}}(\mu_{\bullet}));$$

$$\mathbb{P}(\eta_{\bullet}^N \approx u_{\bullet}) \sim \exp(-N^d \mathcal{I}_{\text{ZRP}}(u_{\bullet}))$$

- ... and to know when we've found the sharpest rate of decay

Introduction and Problem

- Starting point: Kac's model of the spatially homogeneous Boltzmann equation $\mu_t^N \rightarrow \mu_t$, Zero-range process $\eta_t^N \rightarrow u_t$: microscopic models for the nonlinear PDEs

$$\partial_t \mu_t = Q(\mu_t, \mu_t); \quad \partial_t u_t = \frac{1}{2} \Delta \Phi(u_t).$$

- Seek to characterise the large deviations around these limits:

$$\mathbb{P}(\mu_{\bullet}^N \approx \mu_{\bullet}) \sim \exp(-N \mathcal{I}_{\text{KP}}(\mu_{\bullet}));$$

$$\mathbb{P}(\eta_{\bullet}^N \approx u_{\bullet}) \sim \exp(-N^d \mathcal{I}_{\text{ZRP}}(u_{\bullet}))$$

- ... and to know when we've found the sharpest rate of decay (**i.e. matching upper and lower bounds**).

Introduction and Problem

- Starting point: Kac's model of the spatially homogeneous Boltzmann equation $\mu_t^N \rightarrow \mu_t$, Zero-range process $\eta_t^N \rightarrow u_t$: microscopic models for the nonlinear PDEs

$$\partial_t \mu_t = Q(\mu_t, \mu_t); \quad \partial_t u_t = \frac{1}{2} \Delta \Phi(u_t).$$

- Seek to characterise the large deviations around these limits:

$$\mathbb{P}(\mu_{\bullet}^N \approx \mu_{\bullet}) \sim \exp(-N \mathcal{I}_{\text{KP}}(\mu_{\bullet}));$$

$$\mathbb{P}(\eta_{\bullet}^N \approx u_{\bullet}) \sim \exp(-N^d \mathcal{I}_{\text{ZRP}}(u_{\bullet}))$$

- ... and to know when we've found the sharpest rate of decay (**i.e. matching upper and lower bounds**).
- Relation to aspects of the original PDE?

Introduction, 2

- Often find bounds of the form, on a suitable path space \mathbb{X} ,

$$\limsup_N r_N^{-1} \log \mathbb{P}(X_{\bullet}^N \in \mathcal{A}) \leq -\inf \{\mathcal{I}(u) : u \in \mathcal{A}\}$$

for $\mathcal{A} \subset \mathbb{X}$ closed, and

$$\liminf_N r_N^{-1} \log \mathbb{P}(X_{\bullet}^N \in \mathcal{U}) \geq -\inf \{\mathcal{I}(u) : u \in \mathcal{U} \cap \mathcal{X}\}$$

for $\mathcal{U} \subset \mathbb{X}$ open,

Introduction, 2

- Often find bounds of the form, on a suitable path space \mathbb{X} ,

$$\limsup_N r_N^{-1} \log \mathbb{P}(X_{\bullet}^N \in \mathcal{A}) \leq -\inf \{\mathcal{I}(u) : u \in \mathcal{A}\}$$

for $\mathcal{A} \subset \mathbb{X}$ closed, and

$$\liminf_N r_N^{-1} \log \mathbb{P}(X_{\bullet}^N \in \mathcal{U}) \geq -\inf \{\mathcal{I}(u) : u \in \mathcal{U} \cap \mathcal{X}\}$$

for $\mathcal{U} \subset \mathbb{X}$ open, and some universal set $\mathcal{X} \subset \mathbb{X}$.

Introduction, 2

- Often find bounds of the form, on a suitable path space \mathbb{X} ,

$$\limsup_N r_N^{-1} \log \mathbb{P}(X_{\bullet}^N \in \mathcal{A}) \leq -\inf \{\mathcal{I}(u) : u \in \mathcal{A}\}$$

for $\mathcal{A} \subset \mathbb{X}$ closed, and

$$\liminf_N r_N^{-1} \log \mathbb{P}(X_{\bullet}^N \in \mathcal{U}) \geq -\inf \{\mathcal{I}(u) : u \in \mathcal{U} \cap \mathcal{X}\}$$

for $\mathcal{U} \subset \mathbb{X}$ open, and some universal set $\mathcal{X} \subset \mathbb{X}$.

- (Kipnis, Landim for ZRP; Rezakanlou for a collisional model).

Introduction, 2

- Often find bounds of the form, on a suitable path space \mathbb{X} ,

$$\limsup_N r_N^{-1} \log \mathbb{P}(X_{\bullet}^N \in \mathcal{A}) \leq -\inf \{\mathcal{I}(u) : u \in \mathcal{A}\}$$

for $\mathcal{A} \subset \mathbb{X}$ closed, and

$$\liminf_N r_N^{-1} \log \mathbb{P}(X_{\bullet}^N \in \mathcal{U}) \geq -\inf \{\mathcal{I}(u) : u \in \mathcal{U} \cap \mathcal{X}\}$$

for $\mathcal{U} \subset \mathbb{X}$ open, and some universal set $\mathcal{X} \subset \mathbb{X}$.

- (Kipnis, Landim for ZRP; Rezakanlou for a collisional model).
- Leaves open the possibility that there is a better upper bound $\mathcal{I}' > \mathcal{I}$.

Introduction, 2

- Often find bounds of the form, on a suitable path space \mathbb{X} ,

$$\limsup_N r_N^{-1} \log \mathbb{P}(X_{\bullet}^N \in \mathcal{A}) \leq -\inf \{\mathcal{I}(u) : u \in \mathcal{A}\}$$

for $\mathcal{A} \subset \mathbb{X}$ closed, and

$$\liminf_N r_N^{-1} \log \mathbb{P}(X_{\bullet}^N \in \mathcal{U}) \geq -\inf \{\mathcal{I}(u) : u \in \mathcal{U} \cap \mathcal{X}\}$$

for $\mathcal{U} \subset \mathbb{X}$ open, and some universal set $\mathcal{X} \subset \mathbb{X}$.

- (Kipnis, Landim for ZRP; Rezakanlou for a collisional model).
- Leaves open the possibility that there is a better upper bound $\mathcal{I}' > \mathcal{I}$. Can we work harder to get **matching** bounds?

Introduction, 2

- Often find bounds of the form, on a suitable path space \mathbb{X} ,

$$\limsup_N r_N^{-1} \log \mathbb{P}(X_{\bullet}^N \in \mathcal{A}) \leq -\inf \{\mathcal{I}(u) : u \in \mathcal{A}\}$$

for $\mathcal{A} \subset \mathbb{X}$ closed, and

$$\liminf_N r_N^{-1} \log \mathbb{P}(X_{\bullet}^N \in \mathcal{U}) \geq -\inf \{\mathcal{I}(u) : u \in \mathcal{U} \cap \mathcal{X}\}$$

for $\mathcal{U} \subset \mathbb{X}$ open, and some universal set $\mathcal{X} \subset \mathbb{X}$.

- (Kipnis, Landim for ZRP; Rezakanlou for a collisional model).
- Leaves open the possibility that there is a better upper bound $\mathcal{I}' > \mathcal{I}$. Can we work harder to get **matching** bounds?
- Key difficulty: analysis of a modified PDE (controlled Boltzmann equation, skeleton equation)

Introduction, 2

- Often find bounds of the form, on a suitable path space \mathbb{X} ,

$$\limsup_N r_N^{-1} \log \mathbb{P}(X_{\bullet}^N \in \mathcal{A}) \leq -\inf \{\mathcal{I}(u) : u \in \mathcal{A}\}$$

for $\mathcal{A} \subset \mathbb{X}$ closed, and

$$\liminf_N r_N^{-1} \log \mathbb{P}(X_{\bullet}^N \in \mathcal{U}) \geq -\inf \{\mathcal{I}(u) : u \in \mathcal{U} \cap \mathcal{X}\}$$

for $\mathcal{U} \subset \mathbb{X}$ open, and some universal set $\mathcal{X} \subset \mathbb{X}$.

- (Kipnis, Landim for ZRP; Rezakanlou for a collisional model).
- Leaves open the possibility that there is a better upper bound $\mathcal{I}' > \mathcal{I}$. Can we work harder to get **matching** bounds?
- Key difficulty: analysis of a modified PDE (controlled Boltzmann equation, skeleton equation) **with qualitatively different properties**.

Restricted Large Deviations: Kac's Process

- Empirical measure μ_t^N of N interacting velocities $V_1(t), \dots, V_N(t)$.

Kac's Process

- Empirical measure μ_t^N of N interacting velocities $V_1(t), \dots, V_N(t)$.
- Energy-conserving collisions: pairs (v, v_*) update to

$$v \mapsto v' = v + ((v - v_*) \cdot \sigma)\sigma; \quad v_* \mapsto v'_* = v_* + ((v_* - v) \cdot \sigma)\sigma.$$

Kac's Process

- Empirical measure μ_t^N of N interacting velocities $V_1(t), \dots, V_N(t)$.
- Energy-conserving collisions: pairs (v, v_*) update to

$$v \mapsto v' = v + ((v - v_*) \cdot \sigma)\sigma; \quad v_* \mapsto v'_* = v_* + ((v_* - v) \cdot \sigma)\sigma.$$

at rate $N^{-1}B(v - v_*, \sigma)d\sigma = N^{-1}|v - v_*|d\sigma, \sigma \in \mathbb{S}^{d-1}$.

Kac's Process

- Empirical measure μ_t^N of N interacting velocities $V_1(t), \dots, V_N(t)$.
- Energy-conserving collisions: pairs (v, v_*) update to

$$v \mapsto v' = v + ((v - v_*) \cdot \sigma)\sigma; \quad v_* \mapsto v'_* = v_* + ((v_* - v) \cdot \sigma)\sigma.$$

at rate $N^{-1}B(v - v_*, \sigma)d\sigma = N^{-1}|v - v_*|d\sigma, \sigma \in \mathbb{S}^{d-1}$.

- Propagation of chaos: Sznitzman, Grünbaum, Mischler-Mouhot.

Kac's Process

- Empirical measure μ_t^N of N interacting velocities $V_1(t), \dots, V_N(t)$.
- Energy-conserving collisions: pairs (v, v_*) update to

$$v \mapsto v' = v + ((v - v_*) \cdot \sigma)\sigma; \quad v_* \mapsto v'_* = v_* + ((v_* - v) \cdot \sigma)\sigma.$$

at rate $N^{-1}B(v - v_*, \sigma)d\sigma = N^{-1}|v - v_*|d\sigma, \sigma \in \mathbb{S}^{d-1}$.

- Propagation of chaos: Sznitzman, Grünbaum, Mischler-Mouhot.
- For LDP, start in N -particle equilibrium

$$V_0^i \sim_{\text{i.i.d.}} \gamma(dv) = \frac{1}{\sqrt{2\pi d}^d} \exp(-|v|^2/2d)dv.$$

Kac's Process

- Empirical measure μ_t^N of N interacting velocities $V_1(t), \dots, V_N(t)$.
- Energy-conserving collisions: pairs (v, v_*) update to

$$v \mapsto v' = v + ((v - v_*) \cdot \sigma)\sigma; \quad v_* \mapsto v'_* = v_* + ((v_* - v) \cdot \sigma)\sigma.$$

at rate $N^{-1}B(v - v_*, \sigma)d\sigma = N^{-1}|v - v_*|d\sigma, \sigma \in \mathbb{S}^{d-1}$.

- Propagation of chaos: Sznitzman, Grünbaum, Mischler-Mouhot.
- For LDP, start in N -particle equilibrium

$$V_0^i \sim_{\text{i.i.d.}} \gamma(dv) = \frac{1}{\sqrt{2\pi d}^d} \exp(-|v|^2/2d)dv.$$

- Sanov:

$$\mathbb{P}(\mu_0^N \approx \mu_0) \sim \exp(-NH(\mu_0|\gamma)).$$

A Candidate Rate Function

A Candidate Rate Function

- Seek joint LDP on the trajectory $\mu_{\bullet}^N = (\mu_t^N)_{0 \leq t \leq T}$ and the empirical flux w^N recording collisions.

A Candidate Rate Function

- Seek joint LDP on the trajectory $\mu_{\bullet}^N = (\mu_t^N)_{0 \leq t \leq T}$ and the empirical flux w^N recording collisions.
- Candidate rate function: Léonard, 1995; Rezakhanlou, 1998; Bouchet, 2020.

A Candidate Rate Function

- Seek joint LDP on the trajectory $\mu_{\bullet}^N = (\mu_t^N)_{0 \leq t \leq T}$ and the empirical flux w^N recording collisions.
- Candidate rate function: Léonard, 1995; Rezakhanlou, 1998; Bouchet, 2020.
- Dynamic cost: set $\mathcal{J}(\mu_{\bullet}, w) = \infty$ unless (μ_{\bullet}, w) is a measure-flux pair

A Candidate Rate Function

- Seek joint LDP on the trajectory $\mu_{\bullet}^N = (\mu_t^N)_{0 \leq t \leq T}$ and the empirical flux w^N recording collisions.
- Candidate rate function: Léonard, 1995; Rezakhanlou, 1998; Bouchet, 2020.
- Dynamic cost: set $\mathcal{J}(\mu_{\bullet}, w) = \infty$ unless (μ_{\bullet}, w) is a measure-flux pair, in which case set

$$\mathcal{J}(\mu_{\bullet}, w) := \text{Ent}(w \mid \underbrace{|v - v_{\star}| \mu_t(dv) \mu_t(dv_{\star}) dt d\sigma}_{=: \bar{m}_{\mu}}).$$

A Candidate Rate Function

- Seek joint LDP on the trajectory $\mu_\bullet^N = (\mu_t^N)_{0 \leq t \leq T}$ and the empirical flux w^N recording collisions.
- Candidate rate function: Léonard, 1995; Rezakhanlou, 1998; Bouchet, 2020.
- Dynamic cost: set $\mathcal{J}(\mu_\bullet, w) = \infty$ unless (μ_\bullet, w) is a measure-flux pair, in which case set

$$\mathcal{J}(\mu_\bullet, w) := \text{Ent}(w \mid \underbrace{|v - v_\star| \mu_t(dv) \mu_t(dv_\star) dt d\sigma}_{=: \bar{m}_\mu}).$$

- If $\mathcal{J} < \infty$, then μ_\bullet solves a modified Boltzmann equation (BE_K),
 $K = \frac{dw}{d\bar{m}_\mu}$.

A Candidate Rate Function

- Seek joint LDP on the trajectory $\mu_\bullet^N = (\mu_t^N)_{0 \leq t \leq T}$ and the empirical flux w^N recording collisions.
- Candidate rate function: Léonard, 1995; Rezakhanlou, 1998; Bouchet, 2020.
- Dynamic cost: set $\mathcal{J}(\mu_\bullet, w) = \infty$ unless (μ_\bullet, w) is a measure-flux pair, in which case set

$$\mathcal{J}(\mu_\bullet, w) := \text{Ent}(w \mid \underbrace{|v - v_\star| \mu_t(dv) \mu_t(dv_\star) dt d\sigma}_{=: \bar{m}_\mu}).$$

- If $\mathcal{J} < \infty$, then μ_\bullet solves a modified Boltzmann equation (BE_K), $K = \frac{dw}{d\bar{m}_\mu}$. If $\mathcal{J} = 0$, we recover (BE).

A Candidate Rate Function

- Seek joint LDP on the trajectory $\mu_\bullet^N = (\mu_t^N)_{0 \leq t \leq T}$ and the empirical flux w^N recording collisions.
- Candidate rate function: Léonard, 1995; Rezakhanlou, 1998; Bouchet, 2020.
- Dynamic cost: set $\mathcal{J}(\mu_\bullet, w) = \infty$ unless (μ_\bullet, w) is a measure-flux pair, in which case set

$$\mathcal{J}(\mu_\bullet, w) := \text{Ent}(w \mid \underbrace{|v - v_\star| \mu_t(dv) \mu_t(dv_\star) dt d\sigma}_{=: \bar{m}_\mu}).$$

- If $\mathcal{J} < \infty$, then μ_\bullet solves a modified Boltzmann equation (BE $_K$), $K = \frac{dw}{d\bar{m}_\mu}$. If $\mathcal{J} = 0$, we recover (BE).
- Total cost $\mathcal{I}(\mu_\bullet, w) = H(\mu_0 | \gamma) + \mathcal{J}(\mu_\bullet, w)$.

Theorem 1: Positive Result

The rate function \mathcal{I} written above captures at least some of the correct large deviations behaviour:

Theorem (H, 2021; see also Basile-Benedetto-Bertini-Orierri, 2021)

- The variables (μ_{\bullet}^N, w^N) are exponentially tight in $\mathbb{D} \times \mathcal{M}(E)$.
- For all $\mathcal{A} \subset \mathbb{D} \times \mathcal{M}(E)$ closed,

$$\limsup \frac{1}{N} \log \mathbb{P}((\mu_{\bullet}^N, w^N) \in \mathcal{A}) \leq - \inf \{ \mathcal{I}(\mu_{\bullet}, w) : (\mu_{\bullet}, w) \in \mathcal{A} \} .$$

(UB)

Theorem 1: Positive Result

The rate function \mathcal{I} written above captures at least some of the correct large deviations behaviour:

Theorem (H, 2021; see also Basile-Benedetto-Bertini-Orierri, 2021)

- The variables (μ_{\bullet}^N, w^N) are exponentially tight in $\mathbb{D} \times \mathcal{M}(E)$.
- For all $\mathcal{A} \subset \mathbb{D} \times \mathcal{M}(E)$ closed,

$$\limsup \frac{1}{N} \log \mathbb{P}((\mu_{\bullet}^N, w^N) \in \mathcal{A}) \leq - \inf \{ \mathcal{I}(\mu_{\bullet}, w) : (\mu_{\bullet}, w) \in \mathcal{A} \} .$$

(UB)

- For all $\mathcal{U} \subset \mathbb{D} \times \mathcal{M}(E)$ open,

$$\liminf \frac{1}{N} \log \mathbb{P}((\mu_{\bullet}^N, w^N) \in \mathcal{U}) \geq - \inf \{ \mathcal{I}(\mu_{\bullet}, w) : (\mu_{\bullet}, w) \in \mathcal{U} \cap \mathcal{X} \}$$

(RLB)

where $\mathcal{X} = \{(\mu_{\bullet}, w) : \langle |v|^2 + |v_{\star}|^2, w \rangle < \infty \}$.

Positive Result: Main ideas

- **Upper bound:** variational formulation

$$\mathcal{I}(\mu_\bullet, w) = \sup \left\{ \Xi(\mu_\bullet, w, \varphi, f, g)_T : \varphi \in C_{b,v}, f \in C_{b,t,v}^{1,0}, g \in C_c(E) \right\}$$

Positive Result: Main ideas

- **Upper bound:** variational formulation

$$\mathcal{I}(\mu_\bullet, w) = \sup \left\{ \Xi(\mu_\bullet, w, \varphi, f, g)_T : \varphi \in C_{b,v}, f \in C_{b,t,v}^{1,0}, g \in C_c(E) \right\}$$

where Ξ is constructed so that, for each N ,

$Z^N = \exp(N\Xi(\mu_\bullet^N, w^N, \varphi, f, g)_t)$ is a mean 1 martingale.

Positive Result: Main ideas

- **Upper bound:** variational formulation

$$\mathcal{I}(\mu_\bullet, w) = \sup \left\{ \Xi(\mu_\bullet, w, \varphi, f, g)_T : \varphi \in C_{b,v}, f \in C_{b,t,v}^{1,0}, g \in C_c(E) \right\}$$

where Ξ is constructed so that, for each N ,

$Z^N = \exp(N\Xi(\mu_\bullet^N, w^N, \varphi, f, g)_t)$ is a mean 1 martingale.

- (UB) follows from a standard martingale argument.

Positive Result: Main ideas

- **Upper bound:** variational formulation

$$\mathcal{I}(\mu_\bullet, w) = \sup \left\{ \Xi(\mu_\bullet, w, \varphi, f, g)_T : \varphi \in C_{b,v}, f \in C_{b,t,v}^{1,0}, g \in C_c(E) \right\}$$

where Ξ is constructed so that, for each N ,

$Z^N = \exp(N\Xi(\mu_\bullet^N, w^N, \varphi, f, g)_t)$ is a mean 1 martingale.

- (UB) follows from a standard martingale argument.
- First step in **lower bound:** for (μ_\bullet, w) with $|v - v_*|K$ bounded and bounded away from 0, we can write down a change of measure making (μ_\bullet, w) the typical trajectory as $N \rightarrow \infty$

Positive Result: Main ideas

- **Upper bound:** variational formulation

$$\mathcal{I}(\mu_\bullet, w) = \sup \left\{ \Xi(\mu_\bullet, w, \varphi, f, g)_T : \varphi \in C_{b,v}, f \in C_{b,t,v}^{1,0}, g \in C_c(E) \right\}$$

where Ξ is constructed so that, for each N ,

$Z^N = \exp(N\Xi(\mu_\bullet^N, w^N, \varphi, f, g)_t)$ is a mean 1 martingale.

- (UB) follows from a standard martingale argument.
- First step in **lower bound:** for (μ_\bullet, w) with $|v - v_*|K$ bounded and bounded away from 0, we can write down a change of measure making (μ_\bullet, w) the typical trajectory as $N \rightarrow \infty$ with

$$\mathbb{Q}^N \left(\left| \frac{1}{N} \log \frac{d\mathbb{Q}^N}{d\mathbb{P}} - \mathcal{I}(\mu_\bullet, w) \right| > \epsilon \right) \rightarrow 0.$$

Positive Result: Main ideas

- **Upper bound:** variational formulation

$$\mathcal{I}(\mu_\bullet, w) = \sup \left\{ \Xi(\mu_\bullet, w, \varphi, f, g)_T : \varphi \in C_{b,v}, f \in C_{b,t,v}^{1,0}, g \in C_c(E) \right\}$$

where Ξ is constructed so that, for each N ,

$Z^N = \exp(N\Xi(\mu_\bullet^N, w^N, \varphi, f, g)_t)$ is a mean 1 martingale.

- (UB) follows from a standard martingale argument.
- First step in **lower bound:** for (μ_\bullet, w) with $|v - v_*|K$ bounded and bounded away from 0, we can write down a change of measure making (μ_\bullet, w) the typical trajectory as $N \rightarrow \infty$ with

$$\mathbb{Q}^N \left(\left| \frac{1}{N} \log \frac{d\mathbb{Q}^N}{d\mathbb{P}} - \mathcal{I}(\mu_\bullet, w) \right| > \epsilon \right) \rightarrow 0.$$

- **Argument exploits uniqueness for (BE_K) .**

Extending the Lower Bound

- Call the set where we have a lower bound \mathcal{X}_0 .

Extending the Lower Bound

- Call the set where we have a lower bound \mathcal{X}_0 . By a diagonal argument, we automatically get a lower bound on the set \mathcal{X}_1 of (μ_\bullet, w) for which there exists a recovery sequence

$$(\mu_\bullet^{(n)}, w^{(n)}) \in \mathcal{X}_0 : (\mu_\bullet^{(n)}, w^{(n)}) \rightarrow (\mu_\bullet, w), \mathcal{I}(\mu_\bullet^{(n)}, w^{(n)}) \rightarrow \mathcal{I}(\mu_\bullet, w).$$

- (i.e. where the L.S.C. envelope $\overline{\mathcal{I}}|_{\mathcal{X}_0}$ coincides with \mathcal{I}).

Extending the Lower Bound

- Call the set where we have a lower bound \mathcal{X}_0 . By a diagonal argument, we automatically get a lower bound on the set \mathcal{X}_1 of (μ_\bullet, w) for which there exists a recovery sequence

$$(\mu_\bullet^{(n)}, w^{(n)}) \in \mathcal{X}_0 : (\mu_\bullet^{(n)}, w^{(n)}) \rightarrow (\mu_\bullet, w), \mathcal{I}(\mu_\bullet^{(n)}, w^{(n)}) \rightarrow \mathcal{I}(\mu_\bullet, w).$$

- (i.e. where the L.S.C. envelope $\overline{\mathcal{I}}|_{\mathcal{X}_0}$ coincides with \mathcal{I}).
- Argument: suppression of collisions in various ‘bad’ regions of collision space, then convolution $\mu_t \mapsto \mathcal{G}_\delta \star \mu_t$.

Extending the Lower Bound

- Call the set where we have a lower bound \mathcal{X}_0 . By a diagonal argument, we automatically get a lower bound on the set \mathcal{X}_1 of (μ_\bullet, w) for which there exists a recovery sequence

$$(\mu_\bullet^{(n)}, w^{(n)}) \in \mathcal{X}_0 : (\mu_\bullet^{(n)}, w^{(n)}) \rightarrow (\mu_\bullet, w), \mathcal{I}(\mu_\bullet^{(n)}, w^{(n)}) \rightarrow \mathcal{I}(\mu_\bullet, w).$$

- (i.e. where the L.S.C. envelope $\overline{\mathcal{I}}|_{\mathcal{X}_0}$ coincides with \mathcal{I}).
- Argument: suppression of collisions in various 'bad' regions of collision space, then convolution $\mu_t \mapsto g_\delta \star \mu_t$.
 - Truncation argument critically uses $\langle 1 + |v|^2 + |v_\star|^2, w \rangle < \infty$.

Theorem 2: A Counterexample

Theorem 2: A Counterexample

- (RLB) has the prototypical form of a restricted lower bound.

Theorem 2: A Counterexample

- (RLB) has the prototypical form of a restricted lower bound.
- Counterexample without restriction:

Theorem 2: A Counterexample

- (RLB) has the prototypical form of a restricted lower bound.
- Counterexample without restriction:

Theorem (H, 2021; see also Basile-Bernadetto-Bertini-Caglioti 2021)

Let $\Theta : [0, T] \rightarrow [0, \infty)$ be a bounded energy profile satisfying certain technical assumptions. Then the set

$$\mathcal{A}_\Theta = \{(\mu_\bullet, w) : \mu_0 = \gamma, \mu_\bullet \text{ solves (BE)}, w = \bar{m}_\mu, \langle |v|^2, \mu_t \rangle = \Theta(t)\}$$

is compact, nonempty, and $\mathcal{I} = 0$ on \mathcal{A}_Θ . For some $\mathcal{V} \supset \mathcal{A}_\Theta$,

$$\liminf \frac{1}{N} \log \mathbb{P}((\mu_\bullet^N, w^N) \in \mathcal{V}) < 0 = -\inf \{\mathcal{I}(\mu_\bullet, w) : (\mu_\bullet, w) \in \mathcal{A}_\Theta\}$$

while

$$\inf_{\mathcal{U} \supset \mathcal{A}_\Theta} \liminf \frac{1}{N} \log \mathbb{P}((\mu_\bullet^N, w^N) \in \mathcal{U}) > -\infty.$$

Theorem 2: A Counterexample

- (RLB) has the prototypical form of a restricted lower bound.
- Counterexample without restriction:

Theorem (H, 2021; see also Basile-Bernadetto-Bertini-Caglioti 2021)

Let $\Theta : [0, T] \rightarrow [0, \infty)$ be a bounded energy profile satisfying certain technical assumptions. Then the set

$$\mathcal{A}_\Theta = \{(\mu_\bullet, w) : \mu_0 = \gamma, \mu_\bullet \text{ solves (BE)}, w = \bar{m}_\mu, \langle |v|^2, \mu_t \rangle = \Theta(t)\}$$

is compact, nonempty, and $\mathcal{I} = 0$ on \mathcal{A}_Θ . For some $\mathcal{V} \supset \mathcal{A}_\Theta$,

$$\liminf \frac{1}{N} \log \mathbb{P}((\mu_\bullet^N, w^N) \in \mathcal{V}) < 0 = -\inf \{\mathcal{I}(\mu_\bullet, w) : (\mu_\bullet, w) \in \mathcal{A}_\Theta\}$$

while

$$\inf_{\mathcal{U} \supset \mathcal{A}_\Theta} \liminf \frac{1}{N} \log \mathbb{P}((\mu_\bullet^N, w^N) \in \mathcal{U}) > -\infty.$$

....so such behaviour cannot be excluded, but the rate function predicts the exponential occurrence wrongly.

Counterexample: (One Possible) Proof

- Easiest case: $\Theta(t) = 1 + \theta 1(t > 0)$, for some $\theta > 0$.

Counterexample: (One Possible) Proof

- Easiest case: $\Theta(t) = 1 + \theta 1(t > 0)$, for some $\theta > 0$.
- The Gaussian satisfies $0 < \sup(z : \int e^{z|v|^2} \gamma(dv)) = d < \infty$.

Counterexample: (One Possible) Proof

- Easiest case: $\Theta(t) = 1 + \theta 1(t > 0)$, for some $\theta > 0$.
- The Gaussian satisfies $0 < \sup(z : \int e^{z|v|^2} \gamma(dv)) = d < \infty$.
- Construct \mathbb{Q}^N under which

$$V_0^i \sim_{\text{i.i.d.}} \exp(\lambda_{M(N)} |v|^2 1(|v| \geq M(N))) \gamma(dv)$$

with $\lambda_{M(N)} < d$ chosen so that $\mathbb{E}[V_0^i] = 1 + \theta$,

Counterexample: (One Possible) Proof

- Easiest case: $\Theta(t) = 1 + \theta 1(t > 0)$, for some $\theta > 0$.
- The Gaussian satisfies $0 < \sup(z : \int e^{z|v|^2} \gamma(dv)) = d < \infty$.
- Construct \mathbb{Q}^N under which

$$V_0^i \sim_{\text{i.i.d.}} \exp(\lambda_{M(N)} |v|^2 1(|v| \geq M(N))) \gamma(dv)$$

with $\lambda_{M(N)} < d$ chosen so that $\mathbb{E}[V_0^i] = 1 + \theta$, and $M(N) \rightarrow \infty$ slowly enough that, for all $\delta > 0$,

$$\mathbb{Q}^N (|\langle |v|^2, \mu_0^N \rangle - (1 + \theta)| > \delta) \rightarrow 0; \quad \mathbb{Q}^N (W(\mu_0^N, \gamma_{M(N)}) > \delta) \rightarrow 0.$$

Counterexample: (One Possible) Proof

- Easiest case: $\Theta(t) = 1 + \theta 1(t > 0)$, for some $\theta > 0$.
- The Gaussian satisfies $0 < \sup(z : \int e^{z|v|^2} \gamma(dv)) = d < \infty$.
- Construct \mathbb{Q}^N under which

$$V_0^i \sim_{\text{i.i.d.}} \exp(\lambda_{M(N)} |v|^2 1(|v| \geq M(N))) \gamma(dv)$$

with $\lambda_{M(N)} < d$ chosen so that $\mathbb{E}[V_0^i] = 1 + \theta$, and $M(N) \rightarrow \infty$ slowly enough that, for all $\delta > 0$,

$$\mathbb{Q}^N (|\langle |v|^2, \mu_0^N \rangle - (1 + \theta)| > \delta) \rightarrow 0; \quad \mathbb{Q}^N (W(\mu_0^N, \gamma_{M(N)}) > \delta) \rightarrow 0.$$

- Because $\lambda_{M(N)} < d$, $\frac{d\mathbb{Q}^N}{d\mathbb{P}} \leq e^{Nd(1+\theta+\epsilon)}$ with high \mathbb{Q}^N -probability.

Counterexample: (One Possible) Proof

- Easiest case: $\Theta(t) = 1 + \theta 1(t > 0)$, for some $\theta > 0$.
- The Gaussian satisfies $0 < \sup(z : \int e^{z|v|^2} \gamma(dv)) = d < \infty$.
- Construct \mathbb{Q}^N under which

$$V_0^i \sim_{\text{i.i.d.}} \exp(\lambda_{M(N)} |v|^2 1(|v| \geq M(N))) \gamma(dv)$$

with $\lambda_{M(N)} < d$ chosen so that $\mathbb{E}[V_0^i] = 1 + \theta$, and $M(N) \rightarrow \infty$ slowly enough that, for all $\delta > 0$,

$$\mathbb{Q}^N (|\langle |v|^2, \mu_0^N \rangle - (1 + \theta)| > \delta) \rightarrow 0; \quad \mathbb{Q}^N (W(\mu_0^N, \gamma_{M(N)}) > \delta) \rightarrow 0.$$

- Because $\lambda_{M(N)} < d$, $\frac{d\mathbb{Q}^N}{d\mathbb{P}} \leq e^{Nd(1+\theta+\epsilon)}$ with high \mathbb{Q}^N -probability.
- **Chaotic, but not entropically chaotic.**

Counterexample: (One Possible) Proof

- Easiest case: $\Theta(t) = 1 + \theta 1(t > 0)$, for some $\theta > 0$.
- The Gaussian satisfies $0 < \sup(z : \int e^{z|v|^2} \gamma(dv)) = d < \infty$.
- Construct \mathbb{Q}^N under which

$$V_0^i \sim_{\text{i.i.d.}} \exp(\lambda_{M(N)} |v|^2 1(|v| \geq M(N))) \gamma(dv)$$

with $\lambda_{M(N)} < d$ chosen so that $\mathbb{E}[V_0^i] = 1 + \theta$, and $M(N) \rightarrow \infty$ slowly enough that, for all $\delta > 0$,

$$\mathbb{Q}^N (|\langle |v|^2, \mu_0^N \rangle - (1 + \theta)| > \delta) \rightarrow 0; \quad \mathbb{Q}^N (W(\mu_0^N, \gamma_{M(N)}) > \delta) \rightarrow 0.$$

- Because $\lambda_{M(N)} < d$, $\frac{d\mathbb{Q}^N}{d\mathbb{P}} \leq e^{Nd(1+\theta+\epsilon)}$ with high \mathbb{Q}^N -probability.
- **Chaotic, but not entropically chaotic.**
- Using martingale arguments, Pozvner inequality..., all distributional limits (μ_\bullet, w) are supported on \mathcal{A}_Θ , so for all open $\mathcal{U} \supset \mathcal{A}$,

$$\liminf_N N^{-1} \log \mathbb{P}((\mu_\bullet^N, w) \in \mathcal{U}) \geq -d(1 + \theta).$$

Counterexample: Proof, 2

Counterexample: Proof, 2

- Why doesn't the expected bound hold on \mathcal{A} ?

Counterexample: Proof, 2

- Why doesn't the expected bound hold on \mathcal{A} ?
- $\mathcal{J} = 0$ on \mathcal{A}_θ because all paths have $w = \bar{m}_\mu$, and $H(\mu_0|\gamma) = 0$.

Counterexample: Proof, 2

- Why doesn't the expected bound hold on \mathcal{A} ?
- $\mathcal{J} = 0$ on \mathcal{A}_Θ because all paths have $w = \bar{m}_\mu$, and $H(\mu_0|\gamma) = 0$.
- For well-chosen $f \in C_c(\mathbb{R}^d)$, $0 \leq f \leq |v|^2$,

$$\mathcal{V} = \left\{ (\mu_\bullet, w) : \int_{T/2}^T \langle f, \mu_t \rangle dt > \frac{T}{2} \left(1 + \frac{\theta}{2} \right) \right\} \supset \mathcal{A}_\Theta.$$

Counterexample: Proof, 2

- Why doesn't the expected bound hold on \mathcal{A} ?
- $\mathcal{J} = 0$ on \mathcal{A}_θ because all paths have $w = \bar{m}_\mu$, and $H(\mu_0|\gamma) = 0$.
- For well-chosen $f \in C_c(\mathbb{R}^d)$, $0 \leq f \leq |v|^2$,

$$\mathcal{V} = \left\{ (\mu_\bullet, w) : \int_{T/2}^T \langle f, \mu_t \rangle dt > \frac{T}{2} \left(1 + \frac{\theta}{2} \right) \right\} \supset \mathcal{A}_\theta.$$

and

$$\liminf_N \frac{1}{N} \log \mathbb{P} \left((\mu_\bullet^N, w^N) \in \mathcal{V} \right) \leq \liminf_N \frac{1}{N} \log \mathbb{P} \left(\langle |v|^2, \mu_0^N \rangle > 1 + \frac{\theta}{2} \right) < 0$$

Counterexample: Proof, 2

- Why doesn't the expected bound hold on \mathcal{A} ?
- $\mathcal{J} = 0$ on \mathcal{A}_θ because all paths have $w = \bar{m}_\mu$, and $H(\mu_0|\gamma) = 0$.
- For well-chosen $f \in C_c(\mathbb{R}^d)$, $0 \leq f \leq |v|^2$,

$$\mathcal{V} = \left\{ (\mu_\bullet, w) : \int_{T/2}^T \langle f, \mu_t \rangle dt > \frac{T}{2} \left(1 + \frac{\theta}{2} \right) \right\} \supset \mathcal{A}_\theta.$$

and

$$\liminf_N \frac{1}{N} \log \mathbb{P} \left((\mu_\bullet^N, w^N) \in \mathcal{V} \right) \leq \liminf_N \frac{1}{N} \log \mathbb{P} \left(\langle |v|^2, \mu_0^N \rangle > 1 + \frac{\theta}{2} \right) < 0$$

using pathwise energy conservation and Cramér's theorem.

Counterexample: Proof, 2

- Why doesn't the expected bound hold on \mathcal{A} ?
- $\mathcal{J} = 0$ on \mathcal{A}_Θ because all paths have $w = \bar{m}_\mu$, and $H(\mu_0|\gamma) = 0$.
- For well-chosen $f \in C_c(\mathbb{R}^d)$, $0 \leq f \leq |v|^2$,

$$\mathcal{V} = \left\{ (\mu_\bullet, w) : \int_{T/2}^T \langle f, \mu_t \rangle dt > \frac{T}{2} \left(1 + \frac{\theta}{2} \right) \right\} \supset \mathcal{A}_\Theta.$$

and

$$\liminf_N \frac{1}{N} \log \mathbb{P} \left((\mu_\bullet^N, w^N) \in \mathcal{V} \right) \leq \liminf_N \frac{1}{N} \log \mathbb{P} \left(\langle |v|^2, \mu_0^N \rangle > 1 + \frac{\theta}{2} \right) < 0$$

using pathwise energy conservation and Cramér's theorem.

- **Theorem of Basile-Benedetto-Bertini-Caglioti:** counterexample by a different construction, improved rate function which is > 0 on \mathcal{A}_Θ .

Counterexample: Proof, 2

- Why doesn't the expected bound hold on \mathcal{A} ?
- $\mathcal{J} = 0$ on \mathcal{A}_Θ because all paths have $w = \bar{m}_\mu$, and $H(\mu_0|\gamma) = 0$.
- For well-chosen $f \in C_c(\mathbb{R}^d)$, $0 \leq f \leq |v|^2$,

$$\mathcal{V} = \left\{ (\mu_\bullet, w) : \int_{T/2}^T \langle f, \mu_t \rangle dt > \frac{T}{2} \left(1 + \frac{\theta}{2} \right) \right\} \supset \mathcal{A}_\Theta.$$

and

$$\liminf_N \frac{1}{N} \log \mathbb{P} \left((\mu_\bullet^N, w^N) \in \mathcal{V} \right) \leq \liminf_N \frac{1}{N} \log \mathbb{P} \left(\langle |v|^2, \mu_0^N \rangle > 1 + \frac{\theta}{2} \right) < 0$$

using pathwise energy conservation and Cramér's theorem.

- **Theorem of Basile-Benedetto-Bertini-Caglioti:** counterexample by a different construction, improved rate function which is > 0 on \mathcal{A}_Θ .
- LDP rate function **still** not correct for other Boltzmann kernels (e.g. cutoff Maxwell Molecules).

Zero-Range Process: An Example with Matching Bounds

Zero-Range Process

Zero-Range Process

- Fix a (nondecreasing) function $g : \mathbb{N} \rightarrow \mathbb{R}$.

Zero-Range Process

- Fix a (nondecreasing) function $g : \mathbb{N} \rightarrow \mathbb{R}$.
 - Place a bin at each site of $\mathbb{T}_N^d := \{0, N^{-1}, \dots, 1 - N^{-1}\}^d$;

Zero-Range Process

- Fix a (nondecreasing) function $g : \mathbb{N} \rightarrow \mathbb{R}$.
 - Place a bin at each site of $\mathbb{T}_N^d := \{0, N^{-1}, \dots, 1 - N^{-1}\}^d$;
 - $\tilde{\eta}^N$ = empirical partial configuration;

Zero-Range Process

- Fix a (nondecreasing) function $g : \mathbb{N} \rightarrow \mathbb{R}$.
 - Place a bin at each site of $\mathbb{T}_N^d := \{0, N^{-1}, \dots, 1 - N^{-1}\}^d$;
 - $\tilde{\eta}^N$ = empirical partical configuration;
 - At rate $g(\eta^N(x))$, a particle jumps from x to a randomly chosen neighbour y .

Zero-Range Process

- Fix a (nondecreasing) function $g : \mathbb{N} \rightarrow \mathbb{R}$.
 - Place a bin at each site of $\mathbb{T}_N^d := \{0, N^{-1}, \dots, 1 - N^{-1}\}^d$;
 - $\tilde{\eta}^N$ = empirical partical configuration;
 - At rate $g(\eta^N(x))$, a particle jumps from x to a randomly chosen neighbour y .
 - Identify $\tilde{\eta}^N \in L_{\geq 0}^1(\mathbb{T}_N^d) \subset L_{\geq 0}^1(\mathbb{T}^d)$ and give path space \mathbb{D} the Skorokhod topology for a metric inducing weak convergence.

Zero-Range Process

- Fix a (nondecreasing) function $g : \mathbb{N} \rightarrow \mathbb{R}$.
 - Place a bin at each site of $\mathbb{T}_N^d := \{0, N^{-1}, \dots, 1 - N^{-1}\}^d$;
 - $\tilde{\eta}^N$ = empirical partical configuration;
 - At rate $g(\eta^N(x))$, a particle jumps from x to a randomly chosen neighbour y .
 - Identify $\tilde{\eta}^N \in L_{\geq 0}^1(\mathbb{T}_N^d) \subset L_{\geq 0}^1(\mathbb{T}^d)$ and give path space \mathbb{D} the Skorokhod topology for a metric inducing weak convergence.
- *Hydrodynamic* scaling: limit of $\eta_t^N(x) := \tilde{\eta}_{N^2 t}^N(x)$

Zero-Range Process

- Fix a (nondecreasing) function $g : \mathbb{N} \rightarrow \mathbb{R}$.
 - Place a bin at each site of $\mathbb{T}_N^d := \{0, N^{-1}, \dots, 1 - N^{-1}\}^d$;
 - $\tilde{\eta}^N$ = empirical partical configuration;
 - At rate $g(\eta^N(x))$, a particle jumps from x to a randomly chosen neighbour y .
 - Identify $\tilde{\eta}^N \in L_{\geq 0}^1(\mathbb{T}_N^d) \subset L_{\geq 0}^1(\mathbb{T}^d)$ and give path space \mathbb{D} the Skorokhod topology for a metric inducing weak convergence.
- *Hydrodynamic* scaling: limit of $\eta_t^N(x) := \tilde{\eta}_{N^2 t}^N(x)$ (e.g. Kipnis-Landim)

Zero-Range Process

- Fix a (nondecreasing) function $g : \mathbb{N} \rightarrow \mathbb{R}$.
 - Place a bin at each site of $\mathbb{T}_N^d := \{0, N^{-1}, \dots, 1 - N^{-1}\}^d$;
 - $\tilde{\eta}^N$ = empirical partial configuration;
 - At rate $g(\eta^N(x))$, a particle jumps from x to a randomly chosen neighbour y .
 - Identify $\tilde{\eta}^N \in L_{\geq 0}^1(\mathbb{T}_N^d) \subset L_{\geq 0}^1(\mathbb{T}^d)$ and give path space \mathbb{D} the Skorokhod topology for a metric inducing weak convergence.
- *Hydrodynamic* scaling: limit of $\eta_t^N(x) := \tilde{\eta}_{N^2 t}^N(x)$ (e.g. Kipnis-Landim) - (nondegenerate) nonlinear parabolic equation

Zero-Range Process

- Fix a (nondecreasing) function $g : \mathbb{N} \rightarrow \mathbb{R}$.
 - Place a bin at each site of $\mathbb{T}_N^d := \{0, N^{-1}, \dots, 1 - N^{-1}\}^d$;
 - $\tilde{\eta}^N$ = empirical partial configuration;
 - At rate $g(\eta^N(x))$, a particle jumps from x to a randomly chosen neighbour y .
 - Identify $\tilde{\eta}^N \in L_{\geq 0}^1(\mathbb{T}_N^d) \subset L_{\geq 0}^1(\mathbb{T}^d)$ and give path space \mathbb{D} the Skorokhod topology for a metric inducing weak convergence.
- *Hydrodynamic* scaling: limit of $\eta_t^N(x) := \tilde{\eta}_{N^2 t}^N(x)$ (e.g. Kipnis-Landim) - (nondegenerate) nonlinear parabolic equation

$$\partial_t u_t = \Delta \Phi(u_t)$$

Zero-Range Process

- Fix a (nondecreasing) function $g : \mathbb{N} \rightarrow \mathbb{R}$.
 - Place a bin at each site of $\mathbb{T}_N^d := \{0, N^{-1}, \dots, 1 - N^{-1}\}^d$;
 - $\tilde{\eta}^N$ = empirical partial configuration;
 - At rate $g(\eta^N(x))$, a particle jumps from x to a randomly chosen neighbour y .
 - Identify $\tilde{\eta}^N \in L_{\geq 0}^1(\mathbb{T}_N^d) \subset L_{\geq 0}^1(\mathbb{T}^d)$ and give path space \mathbb{D} the Skorokhod topology for a metric inducing weak convergence.
- *Hydrodynamic* scaling: limit of $\eta_t^N(x) := \tilde{\eta}_{N^2 t}^N(x)$ (e.g. Kipnis-Landim) - (nondegenerate) nonlinear parabolic equation

$$\partial_t u_t = \Delta \Phi(u_t)$$

for some Φ determined by the jump rate g .

Large Deviations: Beyond Kipnis-Landim

- Large deviation result of Kipnis-Landim:

Large Deviations: Beyond Kipnis-Landim

- Large deviation result of Kipnis-Landim: variational rate function

$$\mathcal{I}_{\text{Var}}$$

Large Deviations: Beyond Kipnis-Landim

- Large deviation result of Kipnis-Landim: variational rate function \mathcal{I}_{Var} , lower bound restricted to a set \mathcal{X} .

Large Deviations: Beyond Kipnis-Landim

- Large deviation result of Kipnis-Landim: variational rate function \mathcal{I}_{Var} , lower bound restricted to a set \mathcal{X} .
- In the lower bound, \mathcal{I}_{Var} coincides with

$$\mathcal{I}_{\text{FP}}(u_\bullet) := \mathcal{I}_0(u_0) + \frac{1}{2} \inf_{H \in C_{t,x}^{1,3}} \left\{ \int_{t,x} \Phi(u) |\nabla H|^2 : \partial_t u_t = \Delta \Phi(u_t) - \nabla \cdot (\Phi(u) \nabla H) \right\}$$

Large Deviations: Beyond Kipnis-Landim

- Large deviation result of Kipnis-Landim: variational rate function \mathcal{I}_{Var} , lower bound restricted to a set \mathcal{X} .
- In the lower bound, \mathcal{I}_{Var} coincides with

$$\mathcal{I}_{\text{FP}}(u_\bullet) := \mathcal{I}_0(u_0) + \frac{1}{2} \inf_{H \in C_{t,x}^{1,3}} \left\{ \int_{t,x} \Phi(u) |\nabla H|^2 : \partial_t u_t = \Delta \Phi(u_t) - \nabla \cdot (\Phi(u) \nabla H) \right\}$$

(and \mathcal{X} is the set of u_\bullet where the infimum is finite).

Large Deviations: Beyond Kipnis-Landim

- Large deviation result of Kipnis-Landim: variational rate function \mathcal{I}_{Var} , lower bound restricted to a set \mathcal{X} .
- In the lower bound, \mathcal{I}_{Var} coincides with

$$\mathcal{I}_{\text{FP}}(u_\bullet) := \mathcal{I}_0(u_0) + \frac{1}{2} \inf_{H \in C_{t,x}^{1,3}} \left\{ \int_{t,x} \Phi(u) |\nabla H|^2 : \partial_t u_t = \Delta \Phi(u_t) - \nabla \cdot (\Phi(u) \nabla H) \right\}$$

(and \mathcal{X} is the set of u_\bullet where the infimum is finite).

- What is the L.S.C. envelope of \mathcal{I}_{FP} ?

Large Deviations: Beyond Kipnis-Landim

- Large deviation result of Kipnis-Landim: variational rate function \mathcal{I}_{Var} , lower bound restricted to a set \mathcal{X} .
- In the lower bound, \mathcal{I}_{Var} coincides with

$$\mathcal{I}_{\text{FP}}(u_\bullet) := \mathcal{I}_0(u_0) + \frac{1}{2} \inf_{H \in C_{t,x}^{1,3}} \left\{ \int_{t,x} \Phi(u) |\nabla H|^2 : \partial_t u_t = \Delta \Phi(u_t) - \nabla \cdot (\Phi(u) \nabla H) \right\}$$

(and \mathcal{X} is the set of u_\bullet where the infimum is finite).

- What is the L.S.C. envelope of \mathcal{I}_{FP} ? How much more do we need to improve the upper bound?

Large Deviations: Beyond Kipnis-Landim

- Large deviation result of Kipnis-Landim: variational rate function \mathcal{I}_{Var} , lower bound restricted to a set \mathcal{X} .
- In the lower bound, \mathcal{I}_{Var} coincides with

$$\mathcal{I}_{\text{FP}}(u_\bullet) := \mathcal{I}_0(u_0) + \frac{1}{2} \inf_{H \in C_{t,x}^{1,3}} \left\{ \int_{t,x} \Phi(u) |\nabla H|^2 : \partial_t u_t = \Delta \Phi(u_t) - \nabla \cdot (\Phi(u) \nabla H) \right\}$$

(and \mathcal{X} is the set of u_\bullet where the infimum is finite).

- What is the L.S.C. envelope of \mathcal{I}_{FP} ? How much more do we need to improve the upper bound?
- **Theorem (Fehrman-Gess, 2019)** Under general hypotheses on Φ , including all porous medium nonlinearities $\Phi(u) = u^\alpha$, $\alpha \geq 1$, the LSC envelope is given by

$$\mathcal{I}(u_\bullet) = \mathcal{I}_0(u_0) + \frac{1}{2} \inf_{g \in L_{t,x}^2} \left\{ \|g\|_{L_{t,x}^2}^2 : \partial_t u_t = \Delta \Phi(u_t) - \nabla \cdot (\Phi(u)^{1/2} g) \right\}.$$

Large Deviations: Beyond Kipnis-Landim

- Large deviation result of Kipnis-Landim: variational rate function \mathcal{I}_{Var} , lower bound restricted to a set \mathcal{X} .
- In the lower bound, \mathcal{I}_{Var} coincides with

$$\mathcal{I}_{\text{FP}}(\mathbf{u}_\bullet) := \mathcal{I}_0(u_0) + \frac{1}{2} \inf_{H \in \mathcal{C}_{t,x}^{1,3}} \left\{ \int_{t,x} \Phi(u) |\nabla H|^2 : \partial_t u_t = \Delta \Phi(u_t) - \nabla \cdot (\Phi(u) \nabla H) \right\}$$

(and \mathcal{X} is the set of \mathbf{u}_\bullet where the infimum is finite).

- What is the L.S.C. envelope of \mathcal{I}_{FP} ? How much more do we need to improve the upper bound?
- **Theorem (Fehrman-Gess, 2019)** Under general hypotheses on Φ , including all porous medium nonlinearities $\Phi(u) = u^\alpha$, $\alpha \geq 1$, the LSC envelope is given by

$$\mathcal{I}(\mathbf{u}_\bullet) = \mathcal{I}_0(u_0) + \frac{1}{2} \inf_{g \in L_{t,x}^2} \left\{ \|g\|_{L_{t,x}^2}^2 : \partial_t u_t = \Delta \Phi(u_t) - \nabla \cdot (\Phi(u)^{1/2} g) \right\}.$$

(Skeleton equation (Sk_g)).

The Skeleton Equation (Fehrman-Gess 2019)

The Skeleton Equation (Fehrman-Gess 2019)

- (Sk_g) is critical in L_x^1 and supercritical in L_x^p for any $p > 1$.

The Skeleton Equation (Fehrman-Gess 2019)

- (Sk_g) is critical in L_x^1 and supercritical in L_x^p for any $p > 1$.
- **Not a perturbation of a parabolic equation!**

The Skeleton Equation (Fehrman-Gess 2019)

- (Sk_g) is critical in L_x^1 and supercritical in L_x^p for any $p > 1$.
- **Not a perturbation of a parabolic equation!** Rough drift: e.g. LeBris-Lions, Karlssen-Risebro-Ohlberger-Chen need $g \in W_{loc}^{1,1}$, $\operatorname{div}(g) \in L^\infty$.

The Skeleton Equation (Fehrman-Gess 2019)

- (Sk_g) is critical in L_x^1 and supercritical in L_x^p for any $p > 1$.
- **Not a perturbation of a parabolic equation!** Rough drift: e.g. LeBris-Lions, Karlsson-Risebro-Ohlberger-Chen need $g \in W_{\text{loc}}^{1,1}$, $\text{div}(g) \in L^\infty$.
- *A priori* estimate: for a suitable entropy \mathcal{H}_Φ

$$\mathcal{H}_\Phi(u_T) + \int_0^T \underbrace{\|\nabla \Phi^{1/2}(u_s)\|_{L_x^2}^2}_{=:\mathcal{D}_\Phi(u_s)} ds \leq \mathcal{H}_\Phi(u_0) + c \|g\|_{L_{t,x}^2}^2$$

The Skeleton Equation (Fehrman-Gess 2019)

- (Sk_g) is critical in L_x^1 and supercritical in L_x^p for any $p > 1$.
- **Not a perturbation of a parabolic equation!** Rough drift: e.g. LeBris-Lions, Karlsson-Risebro-Ohlberger-Chen need $g \in W_{\text{loc}}^{1,1}$, $\text{div}(g) \in L^\infty$.
- *A priori* estimate: for a suitable entropy \mathcal{H}_Φ

$$\mathcal{H}_\Phi(u_T) + \underbrace{\int_0^T \|\nabla \Phi^{1/2}(u_s)\|_{L_x^2}^2 ds}_{=: \mathcal{D}_\Phi(u_s)} \leq \mathcal{H}_\Phi(u_0) + c \|g\|_{L_{t,x}^2}^2$$

- **Theorem:** existence and uniqueness in

$$\mathcal{R} := \left\{ u_\bullet \in L_t^\infty L_x^1 : u \geq 0, \sup \mathcal{H}_\Phi(u_t) < \infty, \int_0^T \mathcal{D}_\Phi(u_s) ds < \infty \right\}.$$

The Skeleton Equation (Fehrman-Gess 2019)

- (Sk_g) is critical in L_x^1 and supercritical in L_x^p for any $p > 1$.
- **Not a perturbation of a parabolic equation!** Rough drift: e.g. LeBris-Lions, Karlsson-Risebro-Ohlberger-Chen need $g \in W_{\text{loc}}^{1,1}$, $\text{div}(g) \in L^\infty$.
- *A priori* estimate: for a suitable entropy \mathcal{H}_Φ

$$\mathcal{H}_\Phi(u_T) + \underbrace{\int_0^T \|\nabla \Phi^{1/2}(u_s)\|_{L_x^2}^2 ds}_{=: \mathcal{D}_\Phi(u_s)} \leq \mathcal{H}_\Phi(u_0) + c \|g\|_{L_{t,x}^2}^2$$

- **Theorem:** existence and uniqueness in

$$\mathcal{R} := \left\{ u_\bullet \in L_t^\infty L_x^1 : u \geq 0, \sup \mathcal{H}_\Phi(u_t) < \infty, \int_0^T \mathcal{D}_\Phi(u_s) ds < \infty \right\}.$$

Key ideas: Renormalised kinetic solutions (generalising DiPerna-Lions, Ambrosio, LeBris-Lions).

The Skeleton Equation (Fehrman-Gess 2019)

- (Sk_g) is critical in L_x^1 and supercritical in L_x^p for any $p > 1$.
- **Not a perturbation of a parabolic equation!** Rough drift: e.g. LeBris-Lions, Karlsson-Risebro-Ohlberger-Chen need $g \in W_{\text{loc}}^{1,1}$, $\text{div}(g) \in L^\infty$.
- *A priori* estimate: for a suitable entropy \mathcal{H}_Φ

$$\mathcal{H}_\Phi(u_T) + \underbrace{\int_0^T \|\nabla \Phi^{1/2}(u_s)\|_{L_x^2}^2 ds}_{=:\mathcal{D}_\Phi(u_s)} \leq \mathcal{H}_\Phi(u_0) + c \|g\|_{L_{t,x}^2}^2$$

- **Theorem:** existence and uniqueness in

$$\mathcal{R} := \left\{ u_\bullet \in L_t^\infty L_x^1 : u \geq 0, \sup \mathcal{H}_\Phi(u_t) < \infty, \int_0^T \mathcal{D}_\Phi(u_s) ds < \infty \right\}.$$

Key ideas: Renormalised kinetic solutions (generalising DiPerna-Lions, Ambrosio, LeBris-Lions). Uniqueness of renormalised solutions via variable doubling argument.

The Skeleton Equation (Fehrman-Gess 2019)

- (Sk_g) is critical in L_x^1 and supercritical in L_x^p for any $p > 1$.
- **Not a perturbation of a parabolic equation!** Rough drift: e.g. LeBris-Lions, Karlssen-Risebro-Ohlberger-Chen need $g \in W_{\text{loc}}^{1,1}$, $\text{div}(g) \in L^\infty$.
- *A priori* estimate: for a suitable entropy \mathcal{H}_Φ

$$\mathcal{H}_\Phi(u_T) + \underbrace{\int_0^T \|\nabla \Phi^{1/2}(u_s)\|_{L_x^2}^2 ds}_{=:\mathcal{D}_\Phi(u_s)} \leq \mathcal{H}_\Phi(u_0) + c \|g\|_{L_{t,x}^2}^2$$

- **Theorem:** existence and uniqueness in

$$\mathcal{R} := \left\{ u_\bullet \in L_t^\infty L_x^1 : u \geq 0, \sup \mathcal{H}_\Phi(u_t) < \infty, \int_0^T \mathcal{D}_\Phi(u_s) ds < \infty \right\}.$$

Key ideas: Renormalised kinetic solutions (generalising DiPerna-Lions, Ambrosio, LeBris-Lions). Uniqueness of renormalised solutions via variable doubling argument.

- In particular, $\mathcal{I} = \infty$ outside of \mathcal{R} .

Theorem 3: Large Deviations for a Zero-Range Process

- Aim: LDPs for a particle system around PME $\partial_t u = \frac{1}{2} \Delta u^\alpha, \alpha \geq 1$.

Theorem 3: Large Deviations for a Zero-Range Process

- Aim: LDPs for a particle system around PME $\partial_t u = \frac{1}{2} \Delta u^\alpha, \alpha \geq 1$.
- Consider ZRP $\tilde{\eta}_t^N$ with jump rate $g(k) = 2dk^\alpha$.

Theorem 3: Large Deviations for a Zero-Range Process

- Aim: LDPs for a particle system around PME $\partial_t u = \frac{1}{2} \Delta u^\alpha, \alpha \geq 1$.
- Consider ZRP $\tilde{\eta}_t^N$ with jump rate $g(k) = 2dk^\alpha$. The corresponding nonlinearity is still nondegenerate!

Theorem 3: Large Deviations for a Zero-Range Process

- Aim: LDPs for a particle system around PME $\partial_t u = \frac{1}{2} \Delta u^\alpha, \alpha \geq 1$.
- Consider ZRP $\tilde{\eta}_t^N$ with jump rate $g(k) = 2dk^\alpha$. The corresponding nonlinearity is still nondegenerate!
- See (PME) as a limit with a further rescaling $\eta_t^N(x) := \chi_N \tilde{\eta}_{\chi_N^{\alpha-1} t}^N(x)$, $\chi_N \rightarrow 0$.

Theorem 3: Large Deviations for a Zero-Range Process

- Aim: LDPs for a particle system around PME $\partial_t u = \frac{1}{2} \Delta u^\alpha, \alpha \geq 1$.
- Consider ZRP $\tilde{\eta}_t^N$ with jump rate $g(k) = 2dk^\alpha$. The corresponding nonlinearity is still nondegenerate!
- See (PME) as a limit with a further rescaling $\eta_t^N(x) := \chi_N \tilde{\eta}_{\chi_N^{-\alpha} t}^N(x)$, $\chi_N \rightarrow 0$. Impose $N^2 \chi_N^{\min(1, \alpha/2)}$ **bounded as $N \rightarrow \infty$** .

Theorem (H.-Gess, 2023)

Let η_0^N be drawn from a local equilibrium $\rho \in C(\mathbb{T}^d, (0, \infty))$. Then we have **matching** large deviations upper and lower bounds with speed N^d / χ_N , and the rate function given by taking $\mathcal{I}_0(u_0) := \alpha H(u_0 | \rho)$ and nonlinearity $\Phi(u) = u^\alpha$.

Theorem 3: Large Deviations for a Zero-Range Process

- Aim: LDPs for a particle system around PME $\partial_t u = \frac{1}{2} \Delta u^\alpha, \alpha \geq 1$.
- Consider ZRP $\tilde{\eta}_t^N$ with jump rate $g(k) = 2dk^\alpha$. The corresponding nonlinearity is still nondegenerate!
- See (PME) as a limit with a further rescaling $\eta_t^N(x) := \chi_N \tilde{\eta}_{\chi_N^{\alpha-1} t}^N(x)$, $\chi_N \rightarrow 0$. Impose $N^2 \chi_N^{\min(1, \alpha/2)}$ **bounded as $N \rightarrow \infty$** .

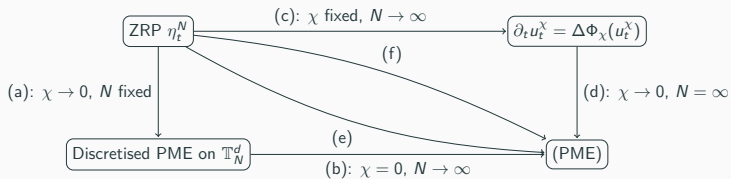
Theorem (H.-Gess, 2023)

Let η_0^N be drawn from a local equilibrium $\rho \in C(\mathbb{T}^d, (0, \infty))$. Then we have **matching** large deviations upper and lower bounds with speed N^d / χ_N , and the rate function given by taking $\mathcal{I}_0(u_0) := \alpha H(u_0 | \rho)$ and nonlinearity $\Phi(u) = u^\alpha$.

- **Innovations: remove paths outside \mathcal{R} by trajectorial estimates, recovery sequences for paths inside \mathcal{R} .**

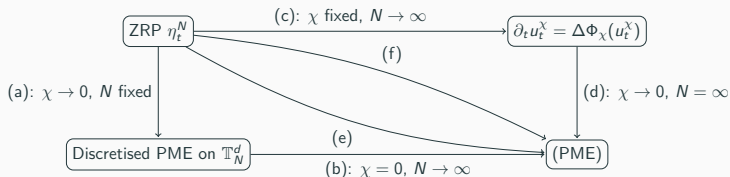
The Scaling Relation

- Possible limits:



The Scaling Relation

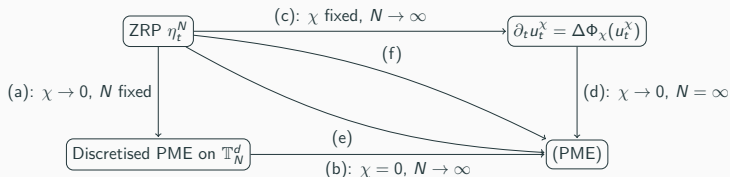
- Possible limits:



- The scaling hypothesis $N^2 \chi_N^{\min(1, \alpha/2)}$ bounded puts us in regime (e).

The Scaling Relation

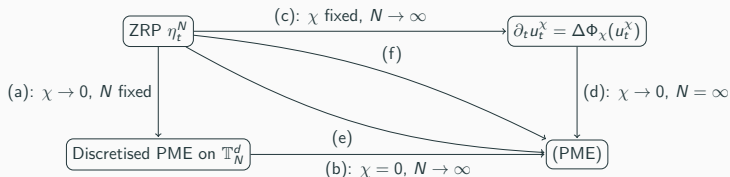
- Possible limits:



- The scaling hypothesis $N^2 \chi_N^{\min(1, \alpha/2)}$ bounded puts us in regime (e).
- (c), (f) are 'classical hydrodynamics';

The Scaling Relation

- Possible limits:

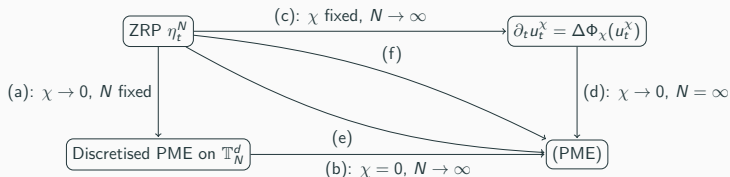


- The scaling hypothesis $N^2 \chi_N^{\min(1, \alpha/2)}$ bounded puts us in regime (e).
- (c), (f) are 'classical hydrodynamics'; (b, e) are *different*.
- Fehrman-Gess: Matching bounds with rate \mathcal{I} for SPDE

$$du_t^\epsilon = \frac{1}{2} \Delta \Phi(u_t^\epsilon) dt - \sqrt{\epsilon} \nabla \cdot (\Phi(u_t)^{1/2} \xi^\delta)$$

The Scaling Relation

- Possible limits:



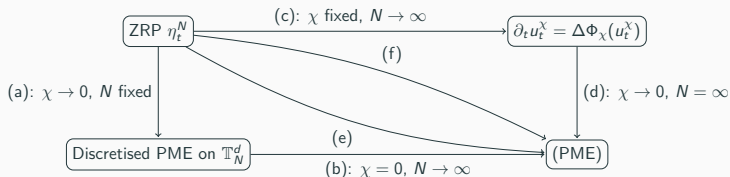
- The scaling hypothesis $N^2 \chi_N^{\min(1, \alpha/2)}$ bounded puts us in regime (e).
- (c), (f) are 'classical hydrodynamics'; (b, e) are *different*.
- Fehrman-Gess: Matching bounds with rate \mathcal{I} for SPDE

$$du_t^\epsilon = \frac{1}{2} \Delta \Phi(u_t^\epsilon) dt - \sqrt{\epsilon} \nabla \cdot (\Phi(u_t)^{1/2} \xi^\delta)$$

with a scaling relation on $(\epsilon, \delta) \rightarrow 0$. Our condition plays the same role.

The Scaling Relation

- Possible limits:



- The scaling hypothesis $N^2 \chi_N^{\min(1, \alpha/2)}$ bounded puts us in regime (e).
- (c), (f) are 'classical hydrodynamics'; (b, e) are *different*.
- Fehrman-Gess: Matching bounds with rate \mathcal{I} for SPDE

$$du_t^\epsilon = \frac{1}{2} \Delta \Phi(u_t^\epsilon) dt - \sqrt{\epsilon} \nabla \cdot (\Phi(u_t)^{1/2} \xi^\delta)$$

with a scaling relation on $(\epsilon, \delta) \rightarrow 0$. Our condition plays the same role.

- 'pathwise regularity in (b,e); rapid local equilibration in (c,f).'

Key Technical Step: Restriction to \mathcal{R} & Replacement Lemma

- Main difficulty in LDP: weak convergence of configuration $\not\Rightarrow$ convergence of nonlinearity (Benois-Kipnis-Landim, Kipnis-Landim).

Key Technical Step: Restriction to \mathcal{R} & Replacement Lemma

- Main difficulty in LDP: weak convergence of configuration $\not\Rightarrow$ convergence of nonlinearity (Benois-Kipnis-Landim, Kipnis-Landim).
- To complete the LDP, we also need to show that, if $u_\bullet \notin \mathcal{R}$ then

$$\inf_{\mathcal{U} \ni u_\bullet} \limsup_N \frac{\chi_N}{N^d} \log \mathbb{P}(\eta_\bullet^N \in \mathcal{U}) = -\infty.$$

Key Technical Step: Restriction to \mathcal{R} & Replacement Lemma

- Main difficulty in LDP: weak convergence of configuration $\not\Rightarrow$ convergence of nonlinearity (Benois-Kipnis-Landim, Kipnis-Landim).
- To complete the LDP, we also need to show that, if $u_\bullet \notin \mathcal{R}$ then

$$\inf_{\mathcal{U} \ni u_\bullet} \limsup_N \frac{\chi_N}{N^d} \log \mathbb{P}(\eta_\bullet^N \in \mathcal{U}) = -\infty.$$

- Kipnis-Landim: replacement lemma in regime (c) via fast macroscopic equilibration.

Key Technical Step: Restriction to \mathcal{R} & Replacement Lemma

- Main difficulty in LDP: weak convergence of configuration $\not\Rightarrow$ convergence of nonlinearity (Benois-Kipnis-Landim, Kipnis-Landim).
- To complete the LDP, we also need to show that, if $u_\bullet \notin \mathcal{R}$ then

$$\inf_{\mathcal{U} \ni u_\bullet} \limsup_N \frac{\chi_N}{N^d} \log \mathbb{P}(\eta_\bullet^N \in \mathcal{U}) = -\infty.$$

- Kipnis-Landim: replacement lemma in regime (c) via fast macroscopic equilibration. For us, both achieved through pathwise regularity using ideas of Aubin-Lions-Simon:

Key Technical Step: Restriction to \mathcal{R} & Replacement Lemma

- Main difficulty in LDP: weak convergence of configuration $\not\Rightarrow$ convergence of nonlinearity (Benois-Kipnis-Landim, Kipnis-Landim).
- To complete the LDP, we also need to show that, if $u_\bullet \notin \mathcal{R}$ then

$$\inf_{\mathcal{U} \ni u_\bullet} \limsup_N \frac{\chi_N}{N^d} \log \mathbb{P}(\eta_\bullet^N \in \mathcal{U}) = -\infty.$$

- Kipnis-Landim: replacement lemma in regime (c) via fast macroscopic equilibration. For us, both achieved through pathwise regularity using ideas of Aubin-Lions-Simon:
 - **Probabilistic Step:** Obtain discrete estimate on $\mathcal{F}_N(\eta_\bullet^N)$ at LDP level

Key Technical Step: Restriction to \mathcal{R} & Replacement Lemma

- Main difficulty in LDP: weak convergence of configuration $\not\Rightarrow$ convergence of nonlinearity (Benois-Kipnis-Landim, Kipnis-Landim).
- To complete the LDP, we also need to show that, if $u_\bullet \notin \mathcal{R}$ then

$$\inf_{\mathcal{U} \ni u_\bullet} \limsup_N \frac{\chi_N}{N^d} \log \mathbb{P}(\eta_\bullet^N \in \mathcal{U}) = -\infty.$$

- Kipnis-Landim: replacement lemma in regime (c) via fast macroscopic equilibration. For us, both achieved through pathwise regularity using ideas of Aubin-Lions-Simon:

- **Probabilistic Step:** Obtain discrete estimate on $\mathcal{F}_N(\eta_\bullet^N)$ at LDP level, for the functional

$$\mathcal{F}_N(\eta_\bullet^N) = \sup_t \mathcal{H}(\eta_t^N) + \frac{1}{N^{d-2}} \int_0^T \sum_{x \sim y} ((\eta_t^N(x))^{\alpha/2} - (\eta_t^N(y))^{\alpha/2})^2 dt.$$

Key Technical Step: Restriction to \mathcal{R} & Replacement Lemma

- Main difficulty in LDP: weak convergence of configuration $\not\Rightarrow$ convergence of nonlinearity (Benois-Kipnis-Landim, Kipnis-Landim).
- To complete the LDP, we also need to show that, if $u_\bullet \notin \mathcal{R}$ then

$$\inf_{\mathcal{U} \ni u_\bullet} \limsup_N \frac{\chi_N}{N^d} \log \mathbb{P}(\eta_\bullet^N \in \mathcal{U}) = -\infty.$$

- Kipnis-Landim: replacement lemma in regime (c) via fast macroscopic equilibration. For us, both achieved through pathwise regularity using ideas of Aubin-Lions-Simon:

- **Probabilistic Step:** Obtain discrete estimate on $\mathcal{F}_N(\eta_\bullet^N)$ at LDP level, for the functional

$$\mathcal{F}_N(\eta_\bullet^N) = \sup_t \mathcal{H}(\eta_t^N) + \frac{1}{N^{d-2}} \int_0^T \sum_{x \sim y} ((\eta_t^N(x))^{\alpha/2} - (\eta_t^N(y))^{\alpha/2})^2 dt.$$

- **Analytic Step:** Pass to the limit on sequences with $\mathcal{F}_N(u_\bullet^N) \leq C$ and $t \mapsto u_t$ continuous;

Key Technical Step: Restriction to \mathcal{R} & Replacement Lemma

- Main difficulty in LDP: weak convergence of configuration $\not\Rightarrow$ convergence of nonlinearity (Benois-Kipnis-Landim, Kipnis-Landim).
- To complete the LDP, we also need to show that, if $u_\bullet \notin \mathcal{R}$ then

$$\inf_{\mathcal{U} \ni u_\bullet} \limsup_N \frac{\chi_N}{N^d} \log \mathbb{P}(\eta_\bullet^N \in \mathcal{U}) = -\infty.$$

- Kipnis-Landim: replacement lemma in regime (c) via fast macroscopic equilibration. For us, both achieved through pathwise regularity using ideas of Aubin-Lions-Simon:

- **Probabilistic Step:** Obtain discrete estimate on $\mathcal{F}_N(\eta_\bullet^N)$ at LDP level, for the functional

$$\mathcal{F}_N(\eta_\bullet^N) = \sup_t \mathcal{H}(\eta_t^N) + \frac{1}{N^{d-2}} \int_0^T \sum_{x \sim y} ((\eta_t^N(x))^{\alpha/2} - (\eta_t^N(y))^{\alpha/2})^2 dt.$$

- **Analytic Step:** Pass to the limit on sequences with $\mathcal{F}_N(u_\bullet^N) \leq C$ and $t \mapsto u_t$ continuous;
- **Probabilistic Step:** eliminate non-continuous u_\bullet

Key Technical Step: Restriction to \mathcal{R} & Replacement Lemma

- Main difficulty in LDP: weak convergence of configuration $\not\Rightarrow$ convergence of nonlinearity (Benois-Kipnis-Landim, Kipnis-Landim).
- To complete the LDP, we also need to show that, if $u_\bullet \notin \mathcal{R}$ then

$$\inf_{\mathcal{U} \ni u_\bullet} \limsup_N \frac{\chi_N}{N^d} \log \mathbb{P}(\eta_\bullet^N \in \mathcal{U}) = -\infty.$$

- Kipnis-Landim: replacement lemma in regime (c) via fast macroscopic equilibration. For us, both achieved through pathwise regularity using ideas of Aubin-Lions-Simon:

- **Probabilistic Step:** Obtain discrete estimate on $\mathcal{F}_N(\eta_\bullet^N)$ at LDP level, for the functional

$$\mathcal{F}_N(\eta_\bullet^N) = \sup_t \mathcal{H}(\eta_t^N) + \frac{1}{N^{d-2}} \int_0^T \sum_{x \sim y} ((\eta_t^N(x))^{\alpha/2} - (\eta_t^N(y))^{\alpha/2})^2 dt.$$

- **Analytic Step:** Pass to the limit on sequences with $\mathcal{F}_N(u_\bullet^N) \leq C$ and $t \mapsto u_t$ continuous;
- **Probabilistic Step:** eliminate non-continuous u_\bullet

PME as Gradient Flow

- PME as gradient flow in space of measures M : Brézis (flat H^{-1}), Otto (Wasserstein).
- New formulation from LDP (Dirr, Peletier, Mielke....)

PME as Gradient Flow

- PME as gradient flow in space of measures \mathcal{M} : Brézis (flat H^{-1}), Otto (Wasserstein).
- New formulation from LDP (Dirr, Peletier, Mielke....): tangent vectors ζ characterised by $\zeta + \nabla \cdot (u^\alpha \nabla \xi) = 0$

PME as Gradient Flow

- PME as gradient flow in space of measures M : Brézis (flat H^{-1}), Otto (Wasserstein).
- New formulation from LDP (Dirr, Peletier, Mielke....): tangent vectors ζ characterised by $\zeta + \nabla \cdot (u^\alpha \nabla \xi) = 0$ with Riemann tensor

$$(\zeta_1, \zeta_2)_u := \int_{\mathbb{T}^d} u^\alpha \nabla \xi_1 \cdot \nabla \xi_2.$$

PME as Gradient Flow

- PME as gradient flow in space of measures M : Brézis (flat H^{-1}), Otto (Wasserstein).
- New formulation from LDP (Dirr, Peletier, Mielke...): tangent vectors ζ characterised by $\zeta + \nabla \cdot (u^\alpha \nabla \xi) = 0$ with Riemann tensor

$$(\zeta_1, \zeta_2)_u := \int_{\mathbb{T}^d} u^\alpha \nabla \xi_1 \cdot \nabla \xi_2.$$

(Generalises Otto's Wasserstein calculus from $\alpha = 1$).

PME as Gradient Flow

- PME as gradient flow in space of measures M : Brézis (flat H^{-1}), Otto (Wasserstein).
- New formulation from LDP (Dirr, Peletier, Mielke...): tangent vectors ζ characterised by $\zeta + \nabla \cdot (u^\alpha \nabla \xi) = 0$ with Riemann tensor

$$(\zeta_1, \zeta_2)_u := \int_{\mathbb{T}^d} u^\alpha \nabla \xi_1 \cdot \nabla \xi_2.$$

(Generalises Otto's Wasserstein calculus from $\alpha = 1$).

- (PME) is formally $\partial_t u = -D_u[\alpha \mathcal{H}(u)]$.

PME as Gradient Flow

- PME as gradient flow in space of measures M : Brézis (flat H^{-1}), Otto (Wasserstein).
- New formulation from LDP (Dirr, Peletier, Mielke...): tangent vectors ζ characterised by $\zeta + \nabla \cdot (u^\alpha \nabla \xi) = 0$ with Riemann tensor

$$(\zeta_1, \zeta_2)_u := \int_{\mathbb{T}^d} u^\alpha \nabla \xi_1 \cdot \nabla \xi_2.$$

(Generalises Otto's Wasserstein calculus from $\alpha = 1$).

- (PME) is formally $\partial_t u = -D_u[\alpha \mathcal{H}(u)]$.
- Manipulate the dynamic cost to the *entropy-dissipation inequality*:

$$\begin{aligned} \mathcal{J}(u_\bullet) &= \frac{1}{2} \int_0^T |\partial_t u + \alpha D\mathcal{H}(u_t)|_{u_t}^2 dt \\ &= \underbrace{\alpha \int_0^T (\partial_t u, D\mathcal{H}(u_t))_{u_t} dt}_{=" \alpha(\mathcal{H}(u_T) - \mathcal{H}(u_0))} + \frac{1}{2} \int_0^T (|\partial_t u_t|_{u_t}^2 + |\alpha D\mathcal{H}(u_t)|_{u_t}^2) dt \end{aligned}$$

PME as Gradient Flow

- PME as gradient flow in space of measures M : Brézis (flat H^{-1}), Otto (Wasserstein).
- New formulation from LDP (Dirr, Peletier, Mielke...): tangent vectors ζ characterised by $\zeta + \nabla \cdot (u^\alpha \nabla \xi) = 0$ with Riemann tensor

$$(\zeta_1, \zeta_2)_u := \int_{\mathbb{T}^d} u^\alpha \nabla \xi_1 \cdot \nabla \xi_2.$$

(Generalises Otto's Wasserstein calculus from $\alpha = 1$).

- (PME) is formally $\partial_t u = -D_u[\alpha \mathcal{H}(u)]$.
- Manipulate the dynamic cost to the *entropy-dissipation inequality*:

$$\begin{aligned} \mathcal{J}(u_\bullet) &= \frac{1}{2} \int_0^T |\partial_t u + \alpha D\mathcal{H}(u_t)|_{u_t}^2 dt \\ &= \underbrace{\alpha \int_0^T (\partial_t u, D\mathcal{H}(u_t))_{u_t} dt}_{=" \alpha(\mathcal{H}(u_T) - \mathcal{H}(u_0))} + \frac{1}{2} \int_0^T (|\partial_t u_t|_{u_t}^2 + |\alpha D\mathcal{H}(u_t)|_{u_t}^2) dt \end{aligned}$$

- Still formally, $|\alpha D\mathcal{H}(u)|_u^2 = \alpha \mathcal{D}_\alpha(u)$.

- Full proof via LDP.

PME as Gradient Flow via LDP

- Full proof via LDP.
- Consider LDP with global equilibrium $\rho = 1$ initial conditions.
Detailed balance $\implies (\mathcal{T}\eta_{\bullet}^N)_t := \eta_{T-t}^N$ has the same law as the original process!

PME as Gradient Flow via LDP

- Full proof via LDP.
- Consider LDP with global equilibrium $\rho = 1$ initial conditions.
Detailed balance $\implies (\mathcal{T}\eta_{\bullet}^N)_t := \eta_{T-t-}^N$ has the same law as the original process!
- Contraction principle and uniqueness of rate functions:

$$\mathcal{I}(\mathcal{T}u_{\bullet}) = \mathcal{I}(u_{\bullet})$$

for all u_{\bullet} .

PME as Gradient Flow via LDP

- Full proof via LDP.
- Consider LDP with global equilibrium $\rho = 1$ initial conditions.
Detailed balance $\implies (\mathcal{T}\eta_{\bullet}^N)_t := \eta_{T-t-}^N$ has the same law as the original process!
- Contraction principle and uniqueness of rate functions:

$$\mathcal{I}(\mathcal{T}u_{\bullet}) = \mathcal{I}(u_{\bullet})$$

for all u_{\bullet} .

- “Improbability of starting at nonequilibrium u_0 and evolving forwards by (PME) = Improbability of evolving via backwards (PME) into u_0 ”.

Theorem 4: Gradient Flow

Theorem (H.-Gess 2023)

Let $u_\bullet \in \mathbb{D}$ with $\mathcal{H}(u_0) < \infty$. Then we have the identity

$$\mathcal{J}(u_\bullet) = \frac{1}{2} \left(\alpha \mathcal{H}(u_T) - \alpha \mathcal{H}(u_0) + \frac{\alpha}{2} \int_0^T \mathcal{D}_\alpha(u_s) ds + \frac{1}{2} \mathcal{A}(u_\bullet) \right) \quad (\text{EDI})$$

allowing both sides to be infinite, where

$$\mathcal{A}(u_\bullet) = \frac{1}{2} \inf \left\{ \|\theta\|_{L^2_{t,x}}^2 : \partial_t u_t + \nabla \cdot \left(\frac{1}{2} u_t^{\alpha/2} \theta_t \right) = 0 \right\}.$$

In particular, the functional on the right-hand side is nonnegative, and vanishes if and only if u_\bullet is a solution to (PME).

Gradient Flow: Sketch Proof

- The unique optimisers for g, θ are characterised by

$$g, \theta \in \Lambda_{u_\bullet} := \overline{\{u^{\alpha/2} \nabla \varphi : \varphi \in C^{1,2}([0, T] \times \mathbb{T}^d)\}}^{L^2_{t,x}}.$$

Gradient Flow: Sketch Proof

- The unique optimisers for g, θ are characterised by

$$g, \theta \in \Lambda_{u_\bullet} := \overline{\{u^{\alpha/2} \nabla \varphi : \varphi \in C^{1,2}([0, T] \times \mathbb{T}^d)\}}^{L^2_{t,x}}.$$

Geometric interpretation: tangent vectors for a.e. t .

Gradient Flow: Sketch Proof

- The unique optimisers for g, θ are characterised by

$$g, \theta \in \Lambda_{u_\bullet} := \overline{\{u^{\alpha/2} \nabla \varphi : \varphi \in C^{1,2}([0, T] \times \mathbb{T}^d)\}}^{L^2_{t,x}}.$$

Geometric interpretation: tangent vectors for a.e. t .

- If g is optimal for u_\bullet , optimal control for $v_\bullet := \mathcal{T}u_\bullet$ is

$$g_r := 2\Pi[v_\bullet] \nabla v^{\alpha/2} - g.$$

Gradient Flow: Sketch Proof

- The unique optimisers for g, θ are characterised by

$$g, \theta \in \Lambda_{u_\bullet} := \overline{\{u^{\alpha/2} \nabla \varphi : \varphi \in C^{1,2}([0, T] \times \mathbb{T}^d)\}}^{L^2_{t,x}}.$$

Geometric interpretation: tangent vectors for a.e. t .

- If g is optimal for u_\bullet , optimal control for $v_\bullet := \mathcal{T}u_\bullet$ is

$$g_r := 2\Pi[v_\bullet] \nabla v^{\alpha/2} - g.$$

Substitute into

$$\mathcal{I}(v_\bullet) = \alpha \mathcal{H}(u_T) + \frac{1}{2} \|g_r\|_{L^2_{t,x}}^2 = \mathcal{I}(u_\bullet) = \alpha \mathcal{H}(u_0) + \frac{1}{2} \|g\|_{L^2_{t,x}}^2.$$

Gradient Flow: Sketch Proof

- The unique optimisers for g, θ are characterised by

$$g, \theta \in \Lambda_{u_\bullet} := \overline{\{u^{\alpha/2} \nabla \varphi : \varphi \in C^{1,2}([0, T] \times \mathbb{T}^d)\}}^{L^2_{t,x}}.$$

Geometric interpretation: tangent vectors for a.e. t .

- If g is optimal for u_\bullet , optimal control for $v_\bullet := \mathcal{T}u_\bullet$ is

$$g_r := 2\Pi[v_\bullet] \nabla v^{\alpha/2} - g.$$

Substitute into

$$\mathcal{I}(v_\bullet) = \alpha \mathcal{H}(u_T) + \frac{1}{2} \|g_r\|_{L^2_{t,x}}^2 = \mathcal{I}(u_\bullet) = \alpha \mathcal{H}(u_0) + \frac{1}{2} \|g\|_{L^2_{t,x}}^2.$$

- After some manipulations,

$$\mathcal{J}(u_\bullet) = \frac{1}{2} \left(\alpha \mathcal{H}(u_T) - \alpha \mathcal{H}(u_0) + \left\| \Pi[u_\bullet] \nabla u^{\alpha/2} \right\|_{L^2_{t,x}}^2 + \mathcal{A}(u_\bullet) \right).$$

Gradient Flow: Sketch Proof, 2

- $\|\Pi[u_\bullet]\nabla u^{\alpha/2}\|_{L^2_{t,x}}^2 \leq \frac{\alpha}{2} \int_0^T \mathcal{D}_\alpha(u_s) ds$, so the previous argument yields the inequality

$$\mathcal{J}(u_\bullet) \leq \frac{1}{2} \left(\alpha \mathcal{H}(u_T) - \alpha \mathcal{H}(u_0) + \frac{\alpha}{2} \int_0^T \mathcal{D}_\alpha(u_s) ds + \frac{1}{2} \mathcal{A}(u_\bullet) \right). \quad (1)$$

Gradient Flow: Sketch Proof, 2

- $\|\Pi[u_\bullet]\nabla u^{\alpha/2}\|_{L^2_{t,x}}^2 \leq \frac{\alpha}{2} \int_0^T \mathcal{D}_\alpha(u_s) ds$, so the previous argument yields the inequality

$$\mathcal{J}(u_\bullet) \leq \frac{1}{2} \left(\alpha \mathcal{H}(u_T) - \alpha \mathcal{H}(u_0) + \frac{\alpha}{2} \int_0^T \mathcal{D}_\alpha(u_s) ds + \frac{1}{2} \mathcal{A}(u_\bullet) \right). \quad (1)$$

- If $u_\bullet \in \mathcal{X}$, then $\nabla u^{\alpha/2} = \frac{2}{\alpha} u^{\alpha/2} \nabla \log u \in \Lambda_{u_\bullet}$, so both of the inequalities are equalities.

Gradient Flow: Sketch Proof, 2

- $\|\Pi[u_\bullet]\nabla u^{\alpha/2}\|_{L^2_{t,x}}^2 \leq \frac{\alpha}{2} \int_0^T \mathcal{D}_\alpha(u_s) ds$, so the previous argument yields the inequality

$$\mathcal{J}(u_\bullet) \leq \frac{1}{2} \left(\alpha \mathcal{H}(u_T) - \alpha \mathcal{H}(u_0) + \frac{\alpha}{2} \int_0^T \mathcal{D}_\alpha(u_s) ds + \frac{1}{2} \mathcal{A}(u_\bullet) \right). \quad (1)$$

- If $u_\bullet \in \mathcal{X}$, then $\nabla u^{\alpha/2} = \frac{2}{\alpha} u^{\alpha/2} \nabla \log u \in \Lambda_{u_\bullet}$, so both of the inequalities are equalities.
- For the general case, use recovery sequences

Gradient Flow: Sketch Proof, 2

- $\|\Pi[u_\bullet]\nabla u^{\alpha/2}\|_{L^2_{t,x}}^2 \leq \frac{\alpha}{2} \int_0^T \mathcal{D}_\alpha(u_s) ds$, so the previous argument yields the inequality

$$\mathcal{J}(u_\bullet) \leq \frac{1}{2} \left(\alpha \mathcal{H}(u_T) - \alpha \mathcal{H}(u_0) + \frac{\alpha}{2} \int_0^T \mathcal{D}_\alpha(u_s) ds + \frac{1}{2} \mathcal{A}(u_\bullet) \right). \quad (1)$$

- If $u_\bullet \in \mathcal{X}$, then $\nabla u^{\alpha/2} = \frac{2}{\alpha} u^{\alpha/2} \nabla \log u \in \Lambda_{u_\bullet}$, so both of the inequalities are equalities.
- For the general case, use recovery sequences and use (1) again.

Gradient Flow: Remark

- LDP* allows us to shortcut proving a 'chain rule for entropy' (Erbar, '16).

Gradient Flow: Remark

- LDP* allows us to shortcut proving a 'chain rule for entropy' (Erbar, '16).
- Same argument: equality in the H -Theorem for (PME)

Gradient Flow: Remark

- LDP* allows us to shortcut proving a ‘chain rule for entropy’ (Erbar, '16).
- Same argument: equality in the H -Theorem for (PME) and (BE), at least for solutions in \mathcal{X}_{BE} .

Gradient Flow: Remark

- LDP* allows us to shortcut proving a 'chain rule for entropy' (Erbar, '16).
- Same argument: equality in the H -Theorem for (PME) and (BE), at least for solutions in \mathcal{X}_{BE} .
- A new look at properties of controlled equations:

Gradient Flow: Remark

- LDP* allows us to shortcut proving a 'chain rule for entropy' (Erbar, '16).
- Same argument: equality in the H -Theorem for (PME) and (BE), at least for solutions in \mathcal{X}_{BE} .
- A new look at properties of controlled equations:
 - Construction of g_r shows how *antidissipative* effects can arise.

Gradient Flow: Remark

- LDP* allows us to shortcut proving a 'chain rule for entropy' (Erbar, '16).
- Same argument: equality in the H -Theorem for (PME) and (BE), at least for solutions in \mathcal{X}_{BE} .
- A new look at properties of controlled equations:
 - Construction of g_r shows how *antidissipative* effects can arise.
 - Hence why L_x^p estimates had to be false

Gradient Flow: Remark

- LDP* allows us to shortcut proving a 'chain rule for entropy' (Erbar, '16).
- Same argument: equality in the H -Theorem for (PME) and (BE), at least for solutions in \mathcal{X}_{BE} .
- A new look at properties of controlled equations:
 - Construction of g_r shows how *antidissipative* effects can arise.
 - Hence why L_x^p estimates had to be false: trajectories with $u_0 \notin L_x^p$, $u_T \in C_x^\infty$ give reversal $v_0 \in C_x^\infty$ but $v_T \notin L_x^p$.

Gradient Flow: Remark

- LDP* allows us to shortcut proving a 'chain rule for entropy' (Erbar, '16).
- Same argument: equality in the H -Theorem for (PME) and (BE), at least for solutions in \mathcal{X}_{BE} .
- A new look at properties of controlled equations:
 - Construction of g_r shows how *antidissipative* effects can arise.
 - Hence why L_x^p estimates had to be false: trajectories with $u_0 \notin L_x^p$, $u_T \in C_x^\infty$ give reversal $v_0 \in C_x^\infty$ but $v_T \notin L_x^p$.
 - Same argument works for (BE $_\kappa$): no possible regularity or moment estimates beyond finite entropy.

Gradient Flow: Remark

- LDP* allows us to shortcut proving a 'chain rule for entropy' (Erbar, '16).
- Same argument: equality in the H -Theorem for (PME) and (BE), at least for solutions in \mathcal{X}_{BE} .
- A new look at properties of controlled equations:
 - Construction of g_r shows how *antidissipative* effects can arise.
 - Hence why L_x^p estimates had to be false: trajectories with $u_0 \notin L_x^p$, $u_T \in C_x^\infty$ give reversal $v_0 \in C_x^\infty$ but $v_T \notin L_x^p$.
 - Same argument works for (BE $_\kappa$): no possible regularity or moment estimates beyond finite entropy.
 - (★): Could be done purely by PDE tools from Fehrman-Gess - but not obvious starting from PME!

- **Kac/Boltzmann:**
 - Heydecker, D., 2023. Large deviations of Kac's conservative particle system and energy nonconserving solutions to the Boltzmann equation: A counterexample to the predicted rate function. *The Annals of Applied Probability*, 33(3), pp.1758-1826.
 - Basile, G., Benedetto, D., Bertini, L. and Caglioti, E., 2022. Asymptotic probability of energy increasing solutions to the homogeneous Boltzmann equation. arXiv preprint arXiv:2202.07311.
- **ZRP / Skeleton Equation:**
 - Gess, B. and Heydecker, D., 2023. A Rescaled Zero-Range Process for the Porous Medium Equation: Hydrodynamic Limit, Large Deviations and Gradient Flow. arXiv preprint arXiv:2303.11289.
 - Fehrman, B. and Gess, B., 2023. Non-equilibrium large deviations and parabolic-hyperbolic PDE with irregular drift. *Inventiones mathematicae*, pp.1-64.