

691 Feb 21 Cayley transform ~~map~~ of some sort between  $U(p, q)$  and  $U(p+q)$ . ~~to~~  
~~the~~ Apparently this is a game played with  $2 \times 2$  matrices, maybe eventually the Gelfand Re-~~gular~~ quasi-determinants - Plücker coordinates. It ~~is~~ amounts to describing a subspace as a graph in two ways. Maybe your ~~observation~~ with length one complexes and general invertible  $2 \times 2$  matrices. 4 objects and 8 maps.

$$\begin{matrix} X_1 \\ \downarrow \\ Y_1 \end{matrix}$$

Let's get it straight  
You want to go between

$$\text{an isom } \begin{matrix} X_1 \\ \oplus \\ Y_1 \end{matrix} \quad \begin{matrix} X_2 \\ \oplus \\ Y_2 \end{matrix}$$

$$Y_1 \quad Y_2 \quad \text{and an isom } \begin{matrix} X_1 \\ \oplus \\ X_2 \end{matrix} \quad \begin{matrix} Y \\ \oplus \\ Y_2 \end{matrix}$$

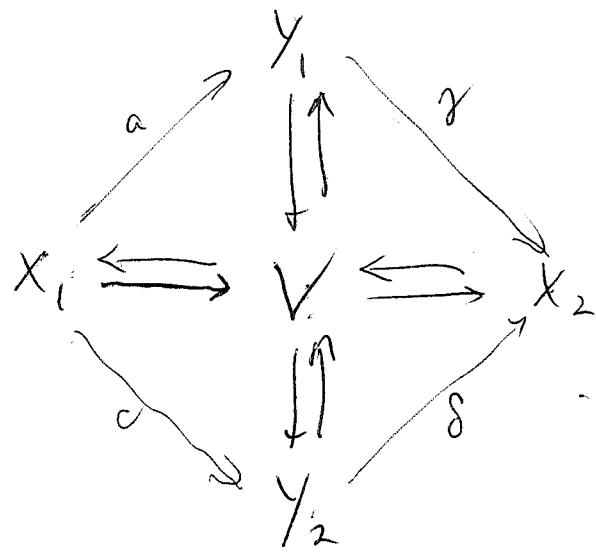
$$\begin{pmatrix} a & b \\ c & d \end{pmatrix} \begin{pmatrix} \alpha & \beta \\ \gamma & \delta \end{pmatrix} = \begin{pmatrix} ax + b\gamma & ab + b\delta \\ cx + d\gamma & c\beta + d\delta \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

$$\begin{pmatrix} \alpha & \beta \\ \gamma & \delta \end{pmatrix} \begin{pmatrix} a & b \\ c & d \end{pmatrix} = \begin{pmatrix} \alpha a + \beta c & \alpha b + \beta d \\ \gamma a + \delta c & \gamma b + \delta d \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

Feb.  
 $\begin{pmatrix} a & b \\ c & d \end{pmatrix} \xrightarrow{\quad Y_1 \quad} \begin{matrix} X_1 \\ \oplus \\ X_2 \end{matrix}$

$$0 \longrightarrow X_1 \longrightarrow X_1 \oplus X_2 \longrightarrow X_2 \longrightarrow 0$$

$$\begin{matrix} & \downarrow \\ & Y_2 \\ & \downarrow \\ & 0 \end{matrix}$$



$$\begin{array}{ccc}
 X_1 & \xleftarrow{\begin{pmatrix} a & b \\ c & d \end{pmatrix}} & Y_1 \\
 \oplus & & \oplus \\
 X_2 & & Y_2
 \end{array}$$

$$\begin{array}{ccccccc}
 & & 0 & & & & \\
 & & \downarrow & & & & \\
 & & Y_1 & \xrightarrow{r} & & & \\
 & & \downarrow f & & & & \\
 0 & \longrightarrow & X_1 & \longrightarrow & V & \longrightarrow & X_2 \longrightarrow 0
 \end{array}$$

A short exact sequences with splitting. But the new point ~~arises~~ arises when one of the maps is an isom.

You start with the sum  $X_1 \oplus X_2 \xrightarrow{\quad} Y_1 \oplus Y_2$

and assume some component is invertible, say  $X_1 \xrightarrow{c} Y_2$ , then you find the "parallel" component  $\delta: Y_1 \rightarrow X_2$  is invertible, explicit inverse given by "quasi-det". Then get ~~map~~  $\boxed{\text{map}}$

$$\begin{array}{ccc}
 X_1 & \xleftarrow{\quad} & X_2 \\
 \oplus & & \oplus \\
 Y_1 & & Y_2
 \end{array}$$

Assume  $c^{-1} \exists$

$$\begin{pmatrix} y_1 \\ y_2 \end{pmatrix} = \begin{pmatrix} a & b \\ c & d \end{pmatrix} \begin{pmatrix} x_1 \\ x_2 \end{pmatrix}$$

$$y_2 = cx_1 + dx_2$$

$$x_1 = c^{-1}(y_2 - dx_2)$$

$$y_1 = ac^{-1}(y_2 - dx_2) + bx_2$$

$$\begin{pmatrix} x_1 \\ y_1 \end{pmatrix} = \begin{pmatrix} -c^{-1}d & c^{-1} \\ b - ac^{-1}d & ac^{-1} \end{pmatrix} \begin{pmatrix} x_2 \\ y_2 \end{pmatrix}$$

$$\begin{array}{ccccccc}
 0 & \longrightarrow & X_1 & \xrightarrow{c} & X & \xrightarrow{} & X_2 \longrightarrow 0 \\
 & & \downarrow & & \downarrow & & \downarrow \\
 0 & \longleftarrow & Y_2 & \longleftarrow & Y & \longleftarrow & Y_1 \longleftarrow 0
 \end{array}$$

693 Point: Given a map  $X \xrightarrow{f} Y$   
 a submodule  $X_1 \subset X$  a quotient module  
 ~~$Y \rightarrow Y/X_1$~~ , such that  ~~$X_1 \subset X \xrightarrow{f} Y \rightarrow Y/X_1$~~  is  
 an isomorphism, then ~~the are canon splittings~~  
 ~~$X = X_1 \oplus f^{-1}(Y_1)$~~   
 ~~$Y = f(X_1) \oplus Y_1$~~   
~~respected by  $f$ .~~  
 So ~~in~~ when one starts with  $X = X_1 \oplus X_2$   
 $Y = Y_1 \oplus Y_2$ , i.e. a choice of complements, you need  
~~to~~ "triangular" matrices to shift to the canonical  
 splittings.

$$\begin{pmatrix} a & b \\ c & d \end{pmatrix} = \begin{pmatrix} & \\ & \end{pmatrix} \begin{pmatrix} & \\ & \end{pmatrix}^*$$

$$\begin{array}{ccccc} Y_2 & f(X_1) & X_1 & X_1 \\ \oplus & \oplus & \oplus & \oplus \\ Y_1 & Y_1 & f^{-1}(Y_1) & X_2 \end{array}$$

What you need now is ~~to add~~ quadratic forms to the situation.

Seems to be the basic ~~mechanism~~ mechanism that you go through when solving equations. So now ~~to~~ see if you can understand the link between  $U(p+q)$  and  $U(p,q)$ . Let us begin with a unitary transformation  $\begin{pmatrix} y_1 \\ y_2 \end{pmatrix} = \begin{pmatrix} a & b \\ c & d \end{pmatrix} \begin{pmatrix} x_1 \\ x_2 \end{pmatrix}$ , and suppose some entry invertible, say  $c$  so that we have a related transf.  $\begin{pmatrix} y_1 \\ x_1 \end{pmatrix} = \begin{pmatrix} a & b \\ d & f \end{pmatrix} \begin{pmatrix} y_2 \\ x_2 \end{pmatrix}$

$$694 \quad \text{then} \quad \begin{aligned} & \text{left out} & & \text{right out} \\ & \|y_1\|^2 + \|y_2\|^2 & = & \|x_1\|^2 + \|x_2\|^2 \\ & \text{left out} & & \text{right in} \\ & \|y_1\|^2 - \|x_1\|^2 & = & \|x_2\|^2 - \|y_2\|^2 \end{aligned}$$

$$\begin{pmatrix} y_1 \\ y_2 \end{pmatrix} = \begin{pmatrix} a & b \\ c & d \end{pmatrix} \begin{pmatrix} x_1 \\ x_2 \end{pmatrix}$$

$$\begin{pmatrix} y_1 \\ x_1 \end{pmatrix} = \begin{pmatrix} \alpha & \beta \\ \gamma & \delta \end{pmatrix} \begin{pmatrix} y_2 \\ x_2 \end{pmatrix}$$

$$x_1 = c^{-1}(y_2 - dx_2)$$

$$y_2 = ac^{-1}(y_2 - dx_2) + bx_2$$

$$\begin{pmatrix} y_1 \\ x_1 \end{pmatrix} = \begin{pmatrix} ac^{-1} & b - ac^{-1}d \\ c^{-1} & -c^{-1}d \end{pmatrix} \begin{pmatrix} y_2 \\ x_2 \end{pmatrix}$$

I'm puzzled by the sign. Consider potential scattering

$$\begin{array}{ccc} x \ll 0 & (-\partial_x^2 + V) u = k^2 u & x \gg 0 \\ \text{left in} \\ e^{+ikx} & \longleftrightarrow & A e^{ikx} + B e^{-ikx} \\ e^{-ikx} & \longleftrightarrow & C e^{ikx} + D e^{-ikx} \end{array}$$

$$\text{Wronskian constant} \Rightarrow \begin{vmatrix} A & B \\ C & D \end{vmatrix} = 1$$

$$\text{reality} \Rightarrow C = \bar{B}, D = \bar{A} \quad \text{when } k \text{ is real}$$

scattering picture  $e^{i(kx - wt)}$  right moving

$$\begin{array}{ccc} \text{left out} & \text{right out} & \text{right in} \\ \frac{C}{D} e^{+ikx} & + e^{-ikx} & \\ \text{left in} \\ e^{ikx} - \frac{B}{D} e^{-ikx} & \longleftrightarrow & \left( A - \frac{BC}{D} \right) e^{ikx} \\ \frac{1}{D} & & \end{array}$$

scattering matrix  $\begin{pmatrix} \frac{C}{D} & \frac{1}{D} \\ \frac{1}{D} & -\frac{B}{D} \end{pmatrix}$

is unitary and complex symmetric

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Start again with a unitary transf.

$$\begin{pmatrix} y_1 \\ y_2 \end{pmatrix} = \begin{pmatrix} a & b \\ c & d \end{pmatrix} \begin{pmatrix} x_1 \\ x_2 \end{pmatrix}$$

$$\|y\|^2 = \|gx\|^2 = \|x\|^2$$

between

$$\begin{matrix} y_1 \\ \oplus \\ y_2 \end{matrix}$$

and

$$\begin{matrix} x_1 \\ \oplus \\ x_2 \end{matrix}$$

Thus

$$x^* g^* g x = x^* x$$

$$\Rightarrow g^* g = I.$$

$$\|y_1\|^2 + \|y_2\|^2 = \|x_1\|^2 + \|x_2\|^2$$

Assume  $c$  invertible whence can solve

$$\begin{pmatrix} y_1 \\ x_1 \end{pmatrix} = \begin{pmatrix} \alpha & \beta \\ \gamma & \delta \end{pmatrix} \begin{pmatrix} y_2 \\ x_2 \end{pmatrix}$$

$$\|y_1\|^2 - \|x_1\|^2 = \|x_2\|^2 - \|y_2\|^2$$

Another possibility

$$\begin{pmatrix} y_1 \\ y_2 \end{pmatrix} = \begin{pmatrix} a & b \\ c & d \end{pmatrix} \begin{pmatrix} x_1 \\ x_2 \end{pmatrix}$$

unitary

$$\|y_1\|^2 + \|y_2\|^2 = \|x_1\|^2 + \|x_2\|^2$$

assume  $a$  invertible

$$\begin{pmatrix} x_1 \\ y_2 \end{pmatrix} = \begin{pmatrix} \alpha & \beta \\ \gamma & \delta \end{pmatrix} \begin{pmatrix} y_1 \\ x_2 \end{pmatrix}$$

$$\|x_1\|^2 - \|y_2\|^2 = \|y_1\|^2 - \|x_2\|^2$$

YES.

Thus the order ~~is~~ important.   Maybe  
the Cramer rule sign applies

696 So where do you go next? You want to understand simple 2 ports. ~~Well good~~ You scattering data  $X \xrightarrow{b} Y$ , ~~dim(X)~~

First case: ~~if~~  $X = 0$ . Then you must give a unitary from ~~V~~  $\xrightarrow{b} W$  between the incoming states and outgoing states. In practice  $V = \mathbb{C}^2 = V_1 \oplus V_2$  and  $W = \mathbb{C}^2 = W_1 \oplus W_2$  are ~~split~~ split, so ~~you~~ you are given a unitary  $\begin{pmatrix} w_1 \\ w_2 \end{pmatrix} = \begin{pmatrix} a & b \\ c & d \end{pmatrix} \begin{pmatrix} v_1 \\ v_2 \end{pmatrix}$  relating incoming and outgoing states. ~~the transmission~~

Goal: You want to ~~be able to~~ be able to handle the graph  $Z \subset R$  with a translation invariant "circuit" with transmission line <sup>segments</sup> linking vertices, a given 2-port used to connect segments at each vertex. ~~idea~~ The idea is then that the port.

The picture: There will be a transfer matrix for the edges. ~~depends on~~ This will be in  $U(1,1)$  depending analytically on  $(z)$

transmission line eqn. ~~for~~

$$\partial_x E + \partial_t I = 0$$

$$\partial_x E + \partial_x E = 0$$

$$(\partial_x + \partial_t)(E + I) = 0$$

$$(\partial_x - \partial_t)(E - I) = 0$$

$$E_x \otimes E_{x+dx} \quad T$$

$$E_x - E_{x+dx} = l dx \frac{\partial}{\partial t} I_x$$

$$I_x - I_{x+dx} = c dx \frac{\partial}{\partial t} E_x$$

$$E + I = A e^{s(-x+t)}$$

$$E - I = B e^{s(x+t)}$$

$$E_0 + I_0 = A$$

$$E_1 + I_1 = A e^{-s}$$

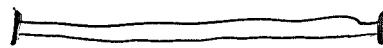
$$E_0 - I_0 = B$$

$$E_1 - I_1 = B e^{-s}$$

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$$\begin{pmatrix} E_0 + I_0 \\ E_0 - I_0 \end{pmatrix} = \begin{pmatrix} e^s & 0 \\ 0 & e^{-s} \end{pmatrix} \begin{pmatrix} E_1 + I_1 \\ E_1 - I_1 \end{pmatrix}$$

So consider a transmission line between  $x=0, 1$ .



$$\xi^+ = E + I \quad \text{fn. of } x$$

$$\xi^- = E - I$$

$$\begin{aligned} (\partial_x + s) \xi^+ &= 0 & \xi^+ &= A e^{-sx} \\ (\partial_x - s) \xi^- &= 0 & \xi^- &= B e^{sx} \end{aligned}$$

$$\boxed{\begin{pmatrix} \xi^+ \\ \xi^- \\ 0 \end{pmatrix} = \begin{pmatrix} e^s & 0 \\ 0 & e^{-s} \end{pmatrix} \begin{pmatrix} \xi^+ \\ \xi^- \\ 0 \end{pmatrix}} \quad \text{transfer matrix}$$

What about the scattering matrix? Incoming coords are  $\xi_0^+, \xi_1^-$  outgoing are  $\xi_0^-, \xi_1^+$

$$\begin{pmatrix} \xi_0^- \\ \xi_1^+ \end{pmatrix} = \begin{pmatrix} e^{-s} & \\ & e^{-s} \end{pmatrix} \begin{pmatrix} \xi_1^- \\ \xi_0^+ \end{pmatrix}$$

Probably better to write

$$\begin{pmatrix} \xi_1^+ \\ \xi_1^- \\ \xi_0^- \end{pmatrix} = \begin{pmatrix} e^{-s} & 0 \\ 0 & e^{-s} \end{pmatrix} \begin{pmatrix} \xi_0^+ \\ \xi_1^- \end{pmatrix}$$

unitary for  $\text{Re}(s) = 0$ .

~~the signal is also decaying in  $\text{Re}(s) > 0$ .~~

698 At ~~the~~<sup>a</sup> vertex you ~~assume~~ first consider  
a 2 point given by a unitary matrix:

$$\text{out} \begin{pmatrix} \xi_0^+ \\ \xi_0^- \\ \xi_1^+ \\ \xi_1^- \end{pmatrix} = \begin{pmatrix} a & b \\ c & d \end{pmatrix} \begin{pmatrix} \xi_0^+ \\ \xi_0^- \\ \xi_1^+ \\ \xi_1^- \end{pmatrix} \quad \text{out}$$

assume  $a^{-1}$  exists, i.e.  $a \neq 0$ , in which case  
from  $|a|^2 + |b|^2 = 1$      $|b|^2 + |d|^2 = 1$ , get  $d \neq 0$   
 $a, d$  are transmission coeffs.  
 $b, c$  — reflection coeffs

$$\begin{pmatrix} \xi_0^+ \\ \xi_0^- \\ \xi_1^+ \\ \xi_1^- \end{pmatrix} = \begin{pmatrix} a^{-1} & -a^{-1}b \\ ca^{-1} & d-ca^{-1}b \end{pmatrix} \begin{pmatrix} \xi_0^+ \\ \xi_0^- \end{pmatrix}$$

$$\xi_1^+ = a \xi_0^+ + b \xi_1^- \quad \xi_0^+ = a^{-1} (\xi_1^+ - b \xi_1^-)$$

$$\xi_0^- = ca^{-1}(\xi_1^+ - b \xi_1^-) + d \xi_1^+$$

~~$\overline{a^{-1}}(-a^{-1}b)$~~   ~~$\overline{ca^{-1}}(d-ca^{-1}b)$~~   
 ~~$+a$~~   
 ~~$|a|^2$~~   ~~$(b + \bar{c})A$~~

~~$ad - ab$~~   
 ~~$a(a+b)$~~

$$\det \begin{pmatrix} a^{-1} & -b \\ c & ad-bc \end{pmatrix} = a^{-2} ad = \frac{d}{a}$$

$$\begin{pmatrix} a & b \\ -b & \bar{a} \end{pmatrix} \rightsquigarrow \begin{pmatrix} a^{-1} & -\frac{b}{a} \\ -\frac{b}{a} & \bar{a} + \frac{|b|^2}{a} \end{pmatrix} = \begin{pmatrix} \frac{1}{a} & -\frac{b}{a} \\ -\frac{b}{a} & \frac{1}{a} \end{pmatrix}$$

699 Rather than proceed with formulae - better

to

$$\begin{pmatrix} a & b \\ -\bar{b} & \bar{a} \end{pmatrix} \mapsto \begin{pmatrix} a^{-1} & -a^{-1}b \\ -\bar{b}(a^{-1}) & \bar{a} + \bar{b}a^{-1}b \end{pmatrix}$$

$$|S| = 1.$$

$$g = \begin{pmatrix} \frac{1}{a} & -\frac{b}{a} \\ -\frac{\bar{b}}{a} & \frac{1}{|a|^2 + |\bar{b}|^2} \end{pmatrix}$$

Suppose  $g = \begin{pmatrix} \alpha & \beta \\ \gamma & \delta \end{pmatrix}$  such that  $g^* \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} g = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$

i.e.,  $\underset{=}{} g^{-1} = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} \alpha^* & \beta^* \\ \gamma^* & \delta^* \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} = \begin{pmatrix} \alpha^* & -\beta^* \\ -\gamma^* & \delta^* \end{pmatrix}$

$$\frac{1}{\Delta} \begin{pmatrix} \delta & -\beta \\ -\gamma & \alpha \end{pmatrix}$$

$$\alpha^* = \frac{\delta}{\Delta}$$

$$\beta^* = \frac{\beta}{\Delta}$$

$$\gamma^* = \frac{\gamma}{\Delta}$$

$$\delta^* = \frac{\alpha}{\Delta}$$

$$\Delta = \det(g) = \frac{1}{a^2} (1 + |b|^2) = g \frac{\bar{a}}{a}$$

~~1~~ ~~2~~

$$\alpha^* = \frac{1}{\bar{a}}$$

$$\frac{\delta}{\Delta} = \frac{\frac{1}{a}}{g \frac{\bar{a}}{a}} = \frac{1}{a}$$

$$\beta^* = -\frac{b}{\bar{a}}$$

$$\frac{\beta}{\Delta} = \left( -\frac{b}{a} \right) \frac{1}{g \frac{\bar{a}}{a}} = -\frac{b}{\bar{a}}$$

$$\gamma^* = -\frac{b}{\bar{a}}$$

$$\frac{\gamma}{\Delta} = \left( -\frac{b}{a} \right) \frac{1}{g \frac{\bar{a}}{a}} = -\frac{b}{\bar{a}}$$

$$\delta^* = \frac{1}{\bar{a}}$$

$$\frac{\alpha}{\Delta} = \frac{1}{a} \frac{1}{g \frac{\bar{a}}{a}} = \frac{1}{\bar{a}} = \frac{1}{\bar{a}}$$

700 So now we have a transfer matrix which is the product of a  $U(1,1)$  matrix (arbitrary) times  $\begin{pmatrix} z & 0 \\ 0 & z^{-1} \end{pmatrix}$ . ~~isogonal field~~

Then what's of interest is the mapping

$$z \mapsto \begin{pmatrix} \alpha & \beta \\ \gamma & \delta \end{pmatrix} \begin{pmatrix} z & 0 \\ 0 & z^{-1} \end{pmatrix} (z)$$

~~and the mapping~~ on the ~~so~~ Riemann sphere.

Actually you should describe determinant 1 :  $SU(1,1)$

$$\begin{pmatrix} \delta - \beta \\ -\gamma \alpha \end{pmatrix} = \begin{pmatrix} z & \beta \\ \gamma & \delta \end{pmatrix}^{-1} = \begin{pmatrix} \alpha^* & -\gamma^* \\ -\beta^* & \delta^* \end{pmatrix} \quad \begin{aligned} \delta &= \alpha^* \\ \beta &= \beta^* \end{aligned}$$

i.e. you have  $\begin{pmatrix} \alpha & \beta \\ \beta^* & \alpha^* \end{pmatrix} \quad |\alpha|^2 - |\beta|^2 = 1.$

We are basically interested in the map on the R.S., Replace  $z$  by  $z^{1/2}$ , and work ~~modulus~~ center

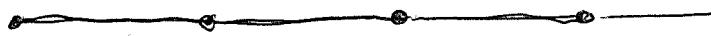
$$\left( \begin{pmatrix} 1 & -h \\ -\bar{h} & 1 \end{pmatrix} \begin{pmatrix} z & 0 \\ 0 & 1 \end{pmatrix} \right) (z) = \frac{z - h}{z + \bar{h}}$$

Simpler: You have  $z \mapsto z \circ$  and then

$$z \mapsto \begin{pmatrix} 1 & -h \\ -\bar{h} & 1 \end{pmatrix} (\circ) = \frac{\circ - h}{1 - \bar{h} \circ} \quad \text{with } |h| < 1.$$

What about the spectrum? ~~There should~~

The point is that the ~~segmented~~ segmented transmission line



has a response which is an analytic function of  $z$ , ~~analytic~~ for  $|z| < 1$  with unitary values for  $|z|=1$ ?

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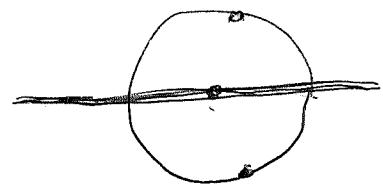
$$S(z) = \begin{pmatrix} 1-h \\ -\bar{h} & 1 \end{pmatrix} (z S(z))$$

$$S = \frac{zS - h}{1 - \bar{h}zS}$$

$$S - \bar{h}zS^2 = zS - h$$

$$0 = \bar{h}zS^2 + (z-1)S - h$$

$$S = \frac{-(z-1) \pm \sqrt{(z-1)^2 + 4\bar{h}zh}}{2\bar{h}z}$$



~~cancel  $(\bar{h}z)^2$~~

$\cos \theta$

$$z^2 - 2z + 1 + 4|h|^2 z$$

$$z^2 - 2z + 1$$

$$= z^2 + 2(2|h|^2 - 1)z + 1$$

range between (-1, 1)

roots:  $z = -(2|h|^2 - 1) \pm \sqrt{(2|h|^2 - 1)^2 - 1}$

$$0 = \bar{h}(zS)^2 + (z^{1/2} - z^{-1/2})(z^{1/2}S) - h$$

Probably can suppose ~~0 < h < 1~~  $0 < h < 1$ .

$$(z^{1/2}S)^2 + \frac{z^{1/2} - z^{-1/2}}{\bar{h}} (z^{1/2}S) - 1$$

Something is funny because  $h=0$   
seems to yield  $S(z) = zS(z)$

702 Feb 22. The idea is to couple transmission line segments in ~~a~~ ~~in~~ a translation invariant way to get some analog of a Schrödinger equation with periodic potential. ~~whose~~ spectrum is continuous but split into bands.

Go over the theory  $(-\partial_x^2 + V(x))\psi = \lambda\psi$

$V(x+2\pi) = V(x)$ . Let  $\Psi(\lambda, x_1, x_0)$  be the propagator

$$(-\partial_x^2 + V(x))\Psi(\lambda, x, x_0) = \lambda\Psi(\lambda, x, x_0)$$

$$\Psi(\lambda, x_0, x_0) = \text{id}.$$

Then  $\Psi(\lambda, 2\pi, 0)$  controls the propagation.

This is a matrix in  $SL_2(\mathbb{R})$  (Wronskian).

~~You ought to be able to construct~~ Point: For each  $\lambda$  you get  $\boxed{\quad}$  2 plane of eigenfunctions. What ~~are the~~ about the

spectrum as ~~a~~ an unbdd s.a. op. on  $L^2(\mathbb{R})$ ? What's important is whether the eigenfunctions are bounded, and that depends on whether the transfer matrix  $\Psi(\lambda, 2\pi, 0)$  is elliptic or hyperbolic. I guess you see instantly that there is no discrete spectrum, just because the product of the eigenvalues is  $\pm 1$ , so ~~either~~ both have abs. value 1, or one has abs. value  $< 1$  and the other  $> 1$ .

$$g^* \varepsilon g = \varepsilon \Rightarrow \overline{\det(g)} \det(\varepsilon) \det(g) = \det(\varepsilon) \\ = |\det(g)| = 1.$$

$$H = H^- \oplus X \oplus H^+$$

$$uH^+ \subset H^+$$

$$u^*H^- \subset H^-$$

$i: X \hookrightarrow H$  inclusion

$i^*i = 1$ .  $\boxed{i^*} i^* = \text{orth. proj. onto } X$ .

Let  $\gamma = i^*ui: X \rightarrow X$ . Claim

$$\begin{cases} i^*u^n i = \gamma^n & n \geq 0 \\ i^*u^{-n} i = (\gamma^*)^n & n \geq 0 \end{cases}$$

Proof:  $u^*H^- \subset H^- \Rightarrow H^- \subset uH^-$

$$u(H^-)^\perp = (uH^-)^\perp \subset (H^-)^\perp, \text{ i.e. } \boxed{\text{H}^- \subset uH^-}$$

$u$  carries  $X \oplus H^+$  into itself, and  $u$  carries  $H^+$  into  $H^+$ , so  $u$  induces an operator  $\gamma$  on  $X \cong (X \oplus H^+)/H^+$ , and  $u^n$  induces  $\gamma^n$  for all  $n \geq 1$ .  $\gamma^n = i^*u^n i$ . Other half:  $(\gamma^*)^n = (i^*u^{-n} i)^* = i^*u^{-n} i$ .

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$$\text{Now look at } H' = \overline{\sum_{n \in \mathbb{Z}} u^n X} \subset H$$

Claim

$$L^2(S, \chi_j d\mu) \xrightarrow{\sim} \boxed{H'} ?$$


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At some point you want to start with a  $\gamma$ .

~~Also~~ First replace  $\gamma$  by  $r\gamma$  where  $0 < r < 1$

so can assume  $\|\gamma\| \leq 1$ , in which case  $1 - \bar{z}\gamma$  etc are invertible.

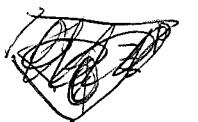
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$$\text{Consider } \mathbb{C}[z, z^{-1}] \otimes_{\mathbb{C}} X = \bigoplus_{n \in \mathbb{Z}} z^n X = H_0$$

with inner product.

$$\begin{aligned} \langle z^p \xi, z^q \xi' \rangle &= \langle \xi, z^{q-p} \xi' \rangle \\ &= \begin{cases} \langle \xi, z^{q-p} \xi' \rangle & q-p \geq 0 \\ \langle \xi, (\gamma^*)^{p-q} \xi' \rangle & q-p \leq 0. \end{cases} \end{aligned}$$

Then



$$\langle \xi, z^n \xi' \rangle_{H_0} = \begin{cases} \langle \xi, \gamma^n \xi' \rangle_X & n \geq 0 \\ \langle \xi, (\gamma^*)^{-n} \xi' \rangle_X & n < 0 \end{cases}$$

$$\langle f(z) \xi, g(z) \xi' \rangle = \langle \xi, f(z)^* g(z) \xi' \rangle$$

$$= \int_{2\pi}^{2\pi} f(z)^* g(z) z^n \langle \xi, (\gamma^*)^n \xi' \rangle$$

$$= \int f(z)^* g(z) \langle \xi, \mu(z) \xi' \rangle$$

$$\mu(z) = \sum_{n \geq 0} z^n \gamma^n + \sum_{n > 0} z^n (\gamma^*)^n$$

$$= \frac{1}{1 - z^{-1} \gamma} + \frac{z \gamma^*}{1 - z \gamma^*}$$

$$= \frac{1}{1 - z \gamma^*} \underbrace{\left( z \gamma^* \left( \frac{1}{1 - z^{-1} \gamma} \right) + \frac{1 - z \gamma^*}{1 - z \gamma^*} \right)}_{1 - \gamma^* g} \frac{1}{1 - z^{-1} \gamma}$$

$$= \frac{1}{1 - z \gamma^*} (1 - \gamma^* g) \frac{1}{1 - z^{-1} \gamma}$$

705 So what happens? You have

$$\begin{aligned} \langle F(z), G(z) \rangle_H &= \int_{2\pi}^{\text{d}\theta} \langle F(z), \frac{1}{(-z)^g} (1-z^*g) \frac{1}{(-z)^f} G(z) \rangle \\ &= \int_{2\pi}^{\text{d}\theta} \left\langle \sqrt{1-z^*g} \frac{1}{(-z)^f} F(z), \sqrt{1-z^*g} \frac{1}{(-z)^g} G(z) \right\rangle \end{aligned}$$

so you find an isometry.

$$\begin{aligned} L^2(S^1, X) &\xrightarrow{\sim} \overline{\bigoplus_{n \in \mathbb{Z}} u^n X} \\ F(z) &\mapsto (1-z^*g)^{1/2} (1-z^{-1}f)^{-1} F(z) \end{aligned}$$

If  $\|f\| < 1$ .

~~so~~

Back to

$$H = H^- \oplus X \oplus H^+$$

$$L^2(S^1, V) \xrightarrow{\sim} \underbrace{\bigoplus_{n \in \mathbb{Z}} u^{-1} V}_{L^2(S^1, X)} \oplus V \oplus X \oplus \underbrace{W \oplus u W \oplus \dots}_{\substack{\oplus \\ \uparrow \\ L^2(S^1, W)}} \longrightarrow L^2(S^1, W)$$

What are the pieces?

$$X \xrightarrow{\frac{a}{b}} Y$$

$$aX \oplus \text{Ker}(a^*)$$

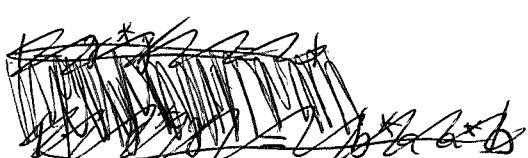
$$|S|$$

$$u = b a^{-1}$$

$$\text{Ker}(b^*) \oplus bX$$

$$\begin{aligned} g &= i^* a i = a^* b a^{-1} a \\ &= a^* b : X \rightarrow X \end{aligned}$$

$$g^* = b^* a : X \rightarrow X.$$



$$\text{Ker}(1-g^*g) ?$$

$$g^*g x = x \Rightarrow \|gx\| = \|x\|.$$

$$\|a^*b x\| = \|bx\| \iff bx \in aX$$

706 So you are confused about  $1 - g^*g$ ?

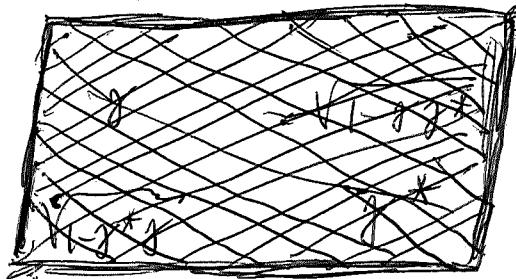
$$\begin{array}{c} ax \oplus w \\ \parallel \\ v' \oplus bw \end{array}$$

~~g = a^\* u c = a^\* b~~

First - dilate a s.a. contraction  $A$ :  $A = A^*$ ,  $A^2 \leq 1$ .  
to an involution.  $F = \begin{pmatrix} A & \sqrt{1-A^2} \\ \sqrt{1-A^2} & -A \end{pmatrix}$

Next - dilate a contraction  $g$  to a unitary

$$\begin{pmatrix} \sqrt{1-g^*g} & g^* \\ g & -\sqrt{1-g^*g} \end{pmatrix} \begin{pmatrix} \sqrt{1-g^*g} & g \\ g^* & -\sqrt{1-g^*g} \end{pmatrix}$$



So the simplest seems to be

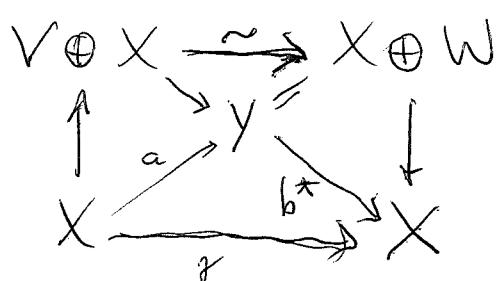
$$\begin{pmatrix} \sqrt{1-g^*g} & g^* \\ g & -\sqrt{1-g^*g} \end{pmatrix}$$

Suppose you have  $X \xleftarrow{g} X'$  a contraction  
and you want to dilate to a unitary, ~~simplest~~  
seems to add

$$\begin{array}{ccc} X & \xleftarrow{g} & X \\ \cancel{\quad \quad \quad} & & \end{array} \quad \text{Yes.}$$

So how does this compare to

$$\begin{array}{ccc} \checkmark & & ax \\ \oplus & & \oplus \\ bX & \xrightarrow{g} & w \end{array}$$



$$= \|(\cancel{a}x_1 + \cancel{b}x_2)\|^2 + \|\sqrt{1-\gamma^*}\gamma x_2\|^2$$

~~$\gamma = b/a$~~

$$\begin{aligned} \|ax_1 + bx_2\|^2 &= \|x_1\|^2 + \|x_2\|^2 + (ax_1, bx_2) + (bx_2, ax_1) \\ &= \|x_1\|^2 + \|x_2\|^2 + (\cancel{x}_1, x_2) + (x_2, \cancel{x}_1) + \|\gamma x_1\|^2 \\ &\quad - \|\gamma x_1\|^2 \\ &= \|\sqrt{1-\gamma^*}\gamma x_1\|^2 + \|\gamma x_1 + x_2\|^2 \end{aligned}$$

$\boxed{\gamma = a^*b}$

$$\|x_1 + \gamma^*x_2\|^2 + \|(1-\gamma\gamma^*)x_2\|^2$$

This corresponds to  $Y = V \oplus X = X \oplus W$

Next  $\|(x_n)\|^2 = \boxed{\dots} \quad \|\sum z^n x_n\|^2 = \sum_{n \leq m} (x_n, \gamma^{m-n} x_m) + \sum_{n > m} (\gamma^{n-m} x_n, x_m)$

~~for all  $x$~~

$$\begin{aligned} \|x_0 + zx_1 + z^2 x_2\|^2 &= \|x_0\|^2 + (\cancel{x}_0, \cancel{\gamma x}_1) + (\cancel{x}_0, \gamma^2 x_2) \\ &\quad + (\cancel{\gamma x}_1, x_0) + \|\gamma x_1\|^2 + (\cancel{\gamma^2 x}_2, x_0) \\ &\quad + (\cancel{\gamma x}_2, x_0) + (\gamma x_1, \cancel{\gamma x}_2) + \|\gamma^2 x_2\|^2 \\ &= u(ax) = bx \\ &\quad + aa^*bx + (1-a^*a)b x \end{aligned}$$

$$= \|x_0\|^2 + (\cancel{\gamma x}_1, x_0) + (\cancel{\gamma^2 x}_2, x_0)$$

$$\begin{array}{ccc} (x_0, \cancel{\gamma x}_1) & \frac{\|x_1\|^2}{\|\gamma x_1\|^2} & (\cancel{\gamma x}_2, x_1) - (\gamma^2 x_2, \cancel{\gamma x}_1) \\ (x_0, \cancel{\gamma^2 x}_2) & (x_1, \cancel{\gamma x}_2) & \frac{\|x_2\|^2}{\|\gamma^2 x_2\|^2} \\ & - (\cancel{\gamma x}_1, \cancel{\gamma^2 x}_2) & - \|\gamma^2 x_2\|^2 \end{array}$$

Recall ~~that~~ complete  $X \oplus X = \{ax_1 + bx_2\}$   
 with  $\|ax_1 + bx_2\|^2 = \|x_1\|^2 + \|x_2\|^2 + \cancel{(x_1, ax_1)} + \cancel{(x_2, bx_2)}$   
 $(x_1, \cancel{a^*b}x_2) + (x_2, \cancel{b^*a}x_1)$   
 $\gamma = b/a$

$$x_0 + ux_1 + u^2x_2 + \dots - (x_0 + vx_1 + v^2x_2 + \dots) \in W \oplus W \oplus u^2W$$

= ? So think. You need a ~~notation~~ notation for ~~a~~  $\tilde{a}^*b$   $\tilde{a}(ix) - i(vx) \in \text{Ker}(a^*) = W$

~~so~~ i.e.

$$\frac{\tilde{a}(ix)}{bx} = \tilde{i}vx + \circled{px}$$

$$aa^*bx + (1-aa^*)bx$$

$$u^2 - v^2 = (v+p)^2 - v^2 = vp + vp + p^2$$

$$= vp + up$$

$$u^2x_2 = v^2x_2 + p(vx_2) + u(px_2)$$

~~so~~ you want to identify  $p$  with  $(1-v^*v)^{1/2}$  ~~so~~  $: X \rightarrow (1-v^*v)^{1/2}X$  eventually.

~~so~~ so think. You have

$$ux = vx + px \in X \oplus W$$

$$u^2x = uvx + upx \quad \cancel{+}$$

$$= v^2x + pvx + upx \in X + W + uW$$

$$u^n x_n = v^n x_n + p v^{n-1} x_n + up v^{n-2} x_n + \dots + u^{n-1} p x_n$$

$$\sum_{n \geq 0} u^n x_n = \underbrace{\sum_{n \geq 0} v^n x_n}_{\in X} + p \underbrace{\sum_{n \geq 1} v^{n-1} x_n}_{\in W} + up \underbrace{\sum_{n \geq 2} v^{n-2} x_n}_{\in W} + \dots$$

$$= \underbrace{\sum_{n \geq 0} v^n x_n}_{\in X} + p \underbrace{\sum_{n \geq 0} v^n x_{n+1}}_{\in W} + up \underbrace{\sum_{n \geq 0} v^n x_{n+2}}_{\in W} + \dots$$

709 ~~with~~ Contraction of  $h: X \rightarrow \cancel{X'}$   
 factor?  $X \xrightarrow{a} Y \xrightarrow{b^*} X'$  where  $a^*a = 1, b^*b = 1.$

$$\text{why } \|ax + bx'\|^2 = \|x\|^2 + \|x'\|^2 + (ax, bx') + (bx', ax) \\ = \|x\|^2 + \|x'\|^2 + (hx, x') + (x', hx)$$

so let  $Y$  be the completion of  $X \oplus X'$  with r.t.  
 this inner product. Have obvious <sup>isom</sup> embeddings  
 $a: X \rightarrow Y, b: X' \rightarrow Y.$

~~Bellairs~~

$$(ax_1 + bx'_1, ax_2 + bx'_2) = (x_1, x_2) + (x'_1, x'_2) + (x'_1, hx_2) + (hx_1, x'_2)$$

How does this construction compare with a Grassmann one

$$\begin{pmatrix} \sqrt{1-h^*h} \\ h \end{pmatrix} X \subset \begin{array}{c} X \\ \oplus \\ X' \end{array} \quad \text{basically the same.}$$

$$a = \begin{pmatrix} \sqrt{1-h^*h} \\ h \end{pmatrix}: X \hookrightarrow \begin{array}{c} X \\ \oplus \\ X' \end{array} \xrightarrow{(0 \ 1)} X'$$

what to do for lecture.

$$H = H^- \oplus X \oplus H^+ \quad \text{o.d.s} \quad u(H^+) \subset H^+ \\ u^{-1}(H^-) \subset H^-$$

Let  $i: X \rightarrow H^\circ$  be the inclusion  
 $i^*: H^- \rightarrow X$  orth proj onto ~~H~~.

$h = i^* u i$  compression of  $u$  to  $X$

$$\text{Claim: } i^* u^n i = \begin{cases} h^n & n \geq 0 \\ (h^*)^{-n} & n \leq 0 \end{cases}$$

Proof.  $u$  carries  $X \oplus H^+ = (H^-)^\perp$  into itself

$H^+$  ~~the quotient space~~

so  $u$  induces an operator  $\bar{u}$  on  $X \oplus H^+ / H^+$ . ~~the quotient space~~  
 namely  $h = i^* u i$ . ~~that's~~ Clearly  $\bar{u}^n = \bar{u}^{-n}$ .

7PD

$$\dim(X) = 1. \quad |h| < 1. \quad \exists \xi_0 \in X \quad |\xi_0| = 1$$

smallest closed subspace of  $H$  containing  $X$  and  
and stable under  $u, u^{-1}$ . Cyclic <sup>unitary</sup> representation of  $\mathbb{Z}$   
There's a Radon measure  $d\mu$

$$L^2(S^1, d\mu) \cong \overline{\bigoplus_{n \in \mathbb{Z}} u^n X} \subset H$$

$$\begin{array}{ccc} z^n \\ f(z) \end{array} \xrightarrow{\hspace{2cm}} \begin{array}{c} u^n \\ f(u) \end{array} \xrightarrow{\hspace{2cm}} \begin{array}{c} \xi_0 \\ \xi_0 \end{array}$$

$$(f, g) = \int \overline{f(z)} g(z) d\mu$$

□  $\int z^n d\mu = (\xi_0, u^n \xi_0) = \begin{cases} h^n & n \geq 0 \\ \bar{h}^{-n} & n < 0. \end{cases}$

$$d\mu = \left( \sum_{n \geq 0} z^n h^n + \sum_{n \geq 0} z^n \bar{h}^{-n} \right) \frac{d\theta}{2\pi}$$

$$\frac{1}{1 - z^1 h} + \frac{z \bar{h}}{1 - z \bar{h}} = \boxed{\text{[Redacted]}}$$

$$= \frac{(-|h|^2)}{(1 - zh)(1 - z^1 h)}$$

$$L^2(S^1, V) \hookrightarrow H \longleftarrow L^2(S^1, W)$$



$$L^2(S, d\mu)$$

7/11 Try again to get a clean lecture.

$h$ -contraction operator  $\approx X \rightarrow X'$

$$a = \begin{pmatrix} \sqrt{1-h^*h} \\ h \end{pmatrix} X \xrightarrow[\oplus]{(1-h^*)} X' \quad b^* a = h$$

$$\|ax + bx'\|^2 = \|x\|^2 + \|x'\|^2 + (ax, bx') + (bx', ax)$$

$$X \oplus X' \quad \parallel$$

$$\begin{aligned} \|x \oplus x'\|^2 &\stackrel{\text{def}}{=} \|x\|^2 + \|x'\|^2 + (\cancel{bx}, x') + (x', \cancel{bx}) \\ &= \|x + b^*x'\|^2 + (x', (1 - \cancel{b}\cancel{b}^*)x') \\ &= \|x' + \cancel{b}(x)\|^2 + (x, (1 - \cancel{b}\cancel{b}^*)x) \end{aligned}$$

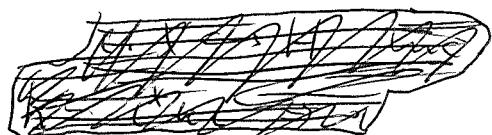
$$Y = X \oplus \underbrace{(1 - \cancel{b}\cancel{b}^*)^{1/2} X'}_{\simeq X'/\text{Ker}(1 - \cancel{b}\cancel{b}^*)^{1/2}}$$

Let's think about what happens over a graph.

~~Sketch~~: February 24

$$H = H^- \oplus X \oplus H^+ \quad u(H^+) \subset H^+ \quad u^{-1}(H^-) \subset H^-$$

$$H = \boxed{\text{graph}} \quad \overbrace{\sum_{n \in \mathbb{Z}} u^n X}^{\text{graph}}$$



$$c : X \hookrightarrow H$$

$$c^* u^n c = \begin{cases} h^n & n \geq 0 \\ (h^*)^{-n} & n \leq 0 \end{cases}$$

712 Coupling: You want to do things over the  $PSL_2(\mathbb{Z})$  tree equivariantly.

For each <sup>valence 3</sup> vertex of the tree we will get a little Hilbert space <sup>X equipped</sup> with contraction acted on by  $\mathbb{Z}/3$ . ~~so~~ We have to divide corresponding  $V, W$  into three pieces permuted cyclically. ~~so~~ For each valence 2 vertex you have another  $(X, h)$  with  $\mathbb{Z}_2$  acting ~~so~~ so splitting of  $V, W$ .

Let's try to make this more explicit.

Each edge has two vertices I think I want ~~one~~ Hilbert space assigned to each edge and <sup>to</sup> each vertex. ~~at each edge~~ At a vertex we have a three port and at an edge we have a 2 port. You now have an understanding of these ports in terms of partial unitaries. A partial unitary ~~will give~~ is:  $X \xrightarrow{\begin{pmatrix} a & \\ b & d+r \end{pmatrix}} Y$ , 3 port means  $r=3$ . Think of  $\text{Ker}(a^*)$  as the "out" gate and  $\text{Ker}(b^*)$  as the "in" gate. ~~at~~ I think I want these gates to split according to incident edges. So what's basically going on is that at each vertex I will get a  $3 \times 3$  unitary operator relating <sup>the</sup> in and out states at the edges and then at each edge we have a  $2 \times 2$  unitary matrix relating in and out states at the two ends.

Anyway, put the structure together

Whatever a port "is" it amounts to something you should be able to pin down.

You want to pin down a port.

Things you are assuming to be true.

1) ~~Given~~ After coupling the components belonging to the edges and vertices you end up with a Hilbert space equipped with a unitary operator. Any  $n$ -port consists of  $n$  gates, each gate gives you an incoming and outgoing line. When you couple two gates you identify the outgoing line of one to the incoming line of the other, and the reverse. This means the Hilbert space "over" ~~the~~ graph consists of the  $\mathcal{Y} = \mathcal{X} \oplus \mathcal{W}$  space for each vertex and for each edge? Wait. Think about connecting partial unitaries to obtain a unitary operator. It seems clear you ~~can~~ can do this easily.

2) You should be able to understand the spectrum of the unitary operator by Gelfand type eigenvectors. By understand I mean seeing whether the spectrum is continuous or discrete. I'm looking for eigenvectors which decay on one side of an edge.

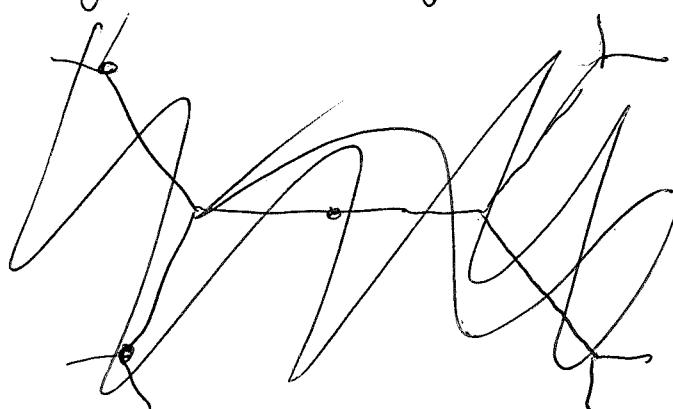
What is the eigenvector equation? Let's see.

You are seeking the Hilbert space. This consists of an  $\mathcal{X}$  space for each vertex

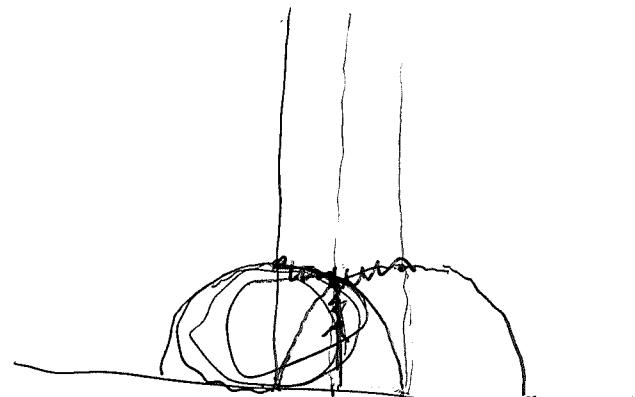
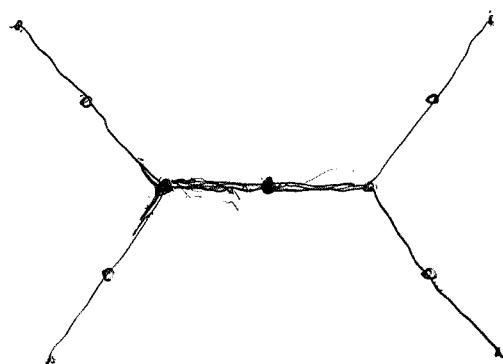
~~Off the Wall~~

You need to organize things. How?

You first have to describe the Hilbert space and the unitary operator. It would be nice if the Hilbert <sup>space</sup> turns out to be some sort of space of  $L^2$  cochains on the graph.



Picture of graph



At each <sup>3 valent</sup> vertex I think of there being a partial unitary  ~~$V \otimes X \rightarrow X \otimes W$~~  with  $\mathbb{Z}/3$  action, and  $V, W$  split into three parts permuted cyclically. So  ~~$\otimes$~~   $V, W$  split according to the edges issuing from  $v$ . ~~It seems we can think at the~~ At a 2-valent vertex you have a partial unitary with  $\mathbb{Z}/2$  action, YES.

Anyway what ~~happened~~ you learned yesterday was to fix the frequency and calculate the corresponding eigenvectors. Strange business. Once you fix a frequency an  $n$ -port becomes a unitary matrix.

Linear graph.

~~Defn~~

2-port consists of 2 gates (gate = polarized 2-dim Hilbert space) and an isomorphism of the underlying pseudo-Hermitian spaces.

Consider two polarized Hilbert spaces which are isomorphic as pseudo-Hilbert spaces. Equivalently a polarized Hilbert space with another polarization. The other polarization is the same as ~~a~~ <sup>strict</sup> contraction operator.

$$H^+ \oplus H^- \quad W = \Gamma_h = \begin{pmatrix} 1 \\ h \end{pmatrix} H^+$$

$$W \quad \text{where } \|hx\| < \|x\| \quad x \neq 0,$$

$$\boxed{\text{Defn}} \quad H^+ \oplus H^- = W \text{ ortho in the } \begin{matrix} \oplus \\ \text{pseudo innerp} \end{matrix} \quad \text{In the end you get}$$

$$W^\perp \quad \left( \begin{array}{cc} \sqrt{1-h^*h} & h^*/\sqrt{1-h^*h} \\ h/\sqrt{1-h^*h} & \sqrt{1-h^*h} \end{array} \right)$$

~~No right do sign~~ ?

So suppose your edge is

Start again. State space at the left end is <sup>of the edge</sup>

~~H~~  $H^+ \oplus H^-$  and at the right end is  $W^+ \oplus W^-$  and the edge at the fixed freq is an isom.  $(\alpha \beta) \begin{pmatrix} \gamma & \delta \end{pmatrix}$  preserving

716 the pseudo-hermitian product:

$$\|\xi^+\|^2 - \|\xi^-\|^2 = \|\eta^+\|^2 - \|\eta^-\|^2$$

which can be ~~be~~ rewritten

$$\|\xi^+\|^2 + \|\eta^-\|^2 = \|\eta^+\|^2 + \|\xi^-\|^2$$

There are confusing aspects here. I have trouble with the signs. ~~Also~~ Also I would like to see the big unitary operator, at least its eigenfunctions.

Look at a 2 port with transmission

What is a two port. It has two ~~gates~~, each gate is a bol. Hilb space  $H^+ \oplus H^-$ ,  $W^+ \oplus W^-$ .  $H^+$   $W^+$  are ~~the space of~~ incoming spaces on the left & right,  $H^-$  and  $W^-$  are the outgoing

~~Phases is not total~~ One gives a bijection between incoming & outgoing states: ~~is~~

$$H^+ \oplus W^+ \cong H^- \oplus W^-$$

Energy is preserved  $\|\xi^+\|^2 + \|\eta^+\|^2 = \|\xi^-\|^2 + \|\eta^-\|^2$ .

~~Amplitudes~~ Can be rewritten

$$\|\xi^+\|^2 - \|\xi^-\|^2 = -\|\eta^+\|^2 + \|\eta^-\|^2$$

net energy in                    net energy out  
at left                            at the right

Assume good transmission: map  $H^+ \rightarrow W^-$  is an isomorphism.

717 We ~~can~~ have a glimpse of cohomology.  
Each edge gives a diagram

$$\begin{array}{ccc} & \swarrow & \searrow \\ H^+ \otimes H^- & & W^+ \oplus W^- \end{array}$$

So at least on the ~~linear~~ linear graph we have a local system, and the only issue is the  $L^2$  character of the cochains, and ~~what happens then~~ this reduces to the eigenvalues of the transfer matrix.

Next look at trivalent graph to see what can be said. At a vertex ~~you~~ have six quantities, parameters, linked by three relations.

~~the 2 port and 3 port~~ Ultimately I should try to describe the state space.

Start from the viewpoint of a Hilbert space with unitary operator. I think it is ~~completely~~ clear that such a pair exists starting from any 2 port with  $\mathbb{Z}/2$  action and any 3 port with  $\mathbb{Z}/3$  action. This Hilbert space would ~~have to arise by~~ arise by placing the 3 port on each vertex and the 2 port on each edge.

$$\textcircled{1} \quad C[G] \otimes_A Y_A = \underset{A}{\textcircled{1}[G]} \otimes (X_A \oplus \overset{W}{W}_A)$$

$$= (\underset{A}{\textcircled{1}[G]} \otimes X_A) \oplus \underset{A}{\textcircled{1}[G]} \otimes W$$

$$\underset{A}{\textcircled{1}[G]} \otimes Y_A = \underset{A}{\textcircled{1}[G]} \otimes (V_A \oplus X_A) = (\underset{A}{\textcircled{1}[G]} \otimes V) \oplus \underset{A}{\textcircled{1}[G]} \otimes X_A$$

$A = \mathbb{Z}/3$ ,  $B = \mathbb{Z}/2$ , A similar partial unitary arises from the edges

$$\begin{aligned} \mathbb{C}[G] \otimes_{\mathcal{B}} Y_B &= \mathbb{C}[G] \otimes_{\mathcal{B}} (X_B \oplus V_B) \\ &= \mathbb{C}[G] \otimes_{\mathcal{B}} X_B \oplus \mathbb{C}[G] \otimes V \end{aligned}$$

and also  $\simeq \mathbb{C}[G] \otimes W \oplus \mathbb{C}[G] \otimes_{\mathcal{B}} X_B$

same  $V, W$  but interpreted with outgoing  
of incoming reversed.  Thus  
I seem to get a nice Hilbert space with unitary. Why?  
You have two partial unitaries which can  
be glued together.

Examples. Consider the case of an iterated 2-port.  
You ~~ought~~ have then? In general you  
ought to ~~do things you~~ describe things  
geometrically. For the  $SL_2(\mathbb{Z})$ -tree this means  
some sort of coefficient ~~the~~ system on the tree.  
What did I describe above?

There are two ~~the~~ types of vertices ~~the~~

~~the~~  $G/A$   $G/B$  and ~~the~~ edges  $G/C$ .

What is the Hilbert space? You give ~~the~~

$Y_A = X_A \oplus W_A \simeq V_A \oplus X_A$  acted on by  $A = \mathbb{Z}/3$

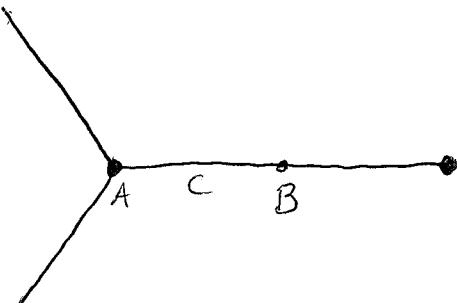
such that  $V_A = \mathbb{C}[A] \otimes_C V_C$ ,  $W_A = \mathbb{C}[A] \otimes_C W_C$ ?  
 $V, W$  associated to the distinguished edge. ~~Similarly~~  
~~the~~ similarly from  $B$  vertex  
give  $Y_B = X_B \oplus (\mathbb{C}[B] \otimes W)$   $= (\mathbb{C}[B] \otimes W) \oplus X_B$

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You ~~should~~ have

$$\mathbb{C}[G] \otimes_A Y_A = \mathbb{C}[G] \otimes_A X_A \oplus \mathbb{C}[G] \otimes_B Y_B$$

I want to work this out very carefully.



$$\text{At } A \text{ we have } Y_A = X_A \oplus \mathbb{C}[A] \otimes_C W_C = \mathbb{C}[A] \otimes_C V_C \oplus X_A$$

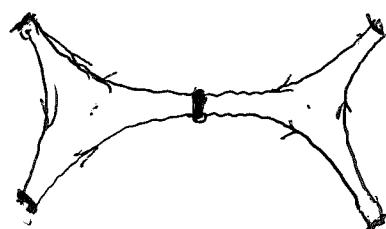
where  $V_C$  = incoming states to A from C $W_C$  = outgoing — from A to C

$$\text{At } B \text{ you have } Y_B = X_B \oplus \quad \text{not correct yet}$$

~~Not~~ Do not try to handle the ~~graph~~

In general you could have  $Y_A, Y_C, Y_B$   
 but instead you want

Feb 26 - ~~Left just like~~ After much fiddling  
 of the tree the good picture seems to be



for handling the situation where the 3 port is ~~is~~ just simple connections appropriate to a ribbon graph. To treat this as a perturbation where the perturbation occurs at the junctions. The free situation is a simple infinite shift for each element of  $P(Q)$ .

720 The first point is the Hilbert space with  $\Gamma = \text{PSL}_2(\mathbb{Z})$  action. There is a basis element for each side of the junction,  $\Gamma$  is acting ~~simply~~ simply transitively on this set. ~~so~~ so we have  $\ell^2(\Gamma)$  with  $\Gamma$  acting on the left and the translation ~~(1)~~  $\begin{pmatrix} 1 & 1 \\ 0 & 1 \end{pmatrix}$  on the right - this is the unperturbed operator. The perturbation takes place at the junction. You want to replace the perfect transmission to allow for some reflection.

Because the operator ~~you are constructing~~ is automatically unitary - Jacobi type etc. ~~There should be a response function associated to one side of the graph and it should satisfy a functional equation.~~

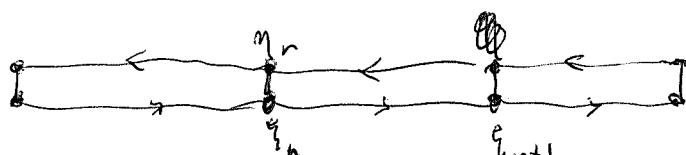
$$f(z) = g(f(z)^2) \quad g \in \mathbb{U}(1,1)$$

and up with a cubic equation? missing  $z$  corresponding to a shift.

Now I have to get the formulas straight.

Begin with the line  $\mathbb{Z}$  acting on  $\mathbb{R}$ .

Picture



Unperturbed situation is  $\ell^2(\mathbb{Z})^2 = L^2(S^1)^2$  with the unitary operator  $\begin{pmatrix} z & 0 \\ 0 & z^{-1} \end{pmatrix}$

721 Thinking of a copy and at each vertex in this picture. Maybe it would help to use  $U, V$  maybe  $U_n, V_n$ . Keep trying.

What you want is two 2-ports connected in alternating fashion  $\alpha \beta \alpha \beta \alpha \beta \dots$ . Let's give an basis  $\xi_n, \eta_n$   $n \in \mathbb{Z}$ . ~~basic perturbation unperturbed~~

$$\begin{cases} u_n \xi_n = \xi_{n+1} \\ u_n \eta_n = \eta_{n-1} \end{cases}$$

perturbed operator will involve transmission and reflector

$$u_n \xi_n = T \xi_{n+1} + \rho \eta_n \quad |\tau|^2 + |\rho|^2 = 1.$$

$$u_n \eta_{n+1} = ? \xi_{n+1} ? \eta_n$$

$$u_n \eta_n = ? \xi_n + ? \eta_{n-1} \quad \text{won't work.}$$

~~Let's start with  $H, \alpha$~~

In the end you want to start with ~~a~~ a transfer matrix in  $U(1,1)$  and produce a unitary operator on  $\ell^2(\mathbb{Z})^{\oplus 2}$  commuting with translation. Picture: The transfer matrix yields a 2-port which can be coupled periodically to itself. (Look up Quantum Hall Effect - Bellessard)

2-port ~~transfer matrix~~ transfer matrix  $\begin{pmatrix} \sqrt{1-\frac{1}{\lambda}} & \frac{1}{\lambda} \\ \frac{1}{\lambda} & \sqrt{1-\frac{1}{\lambda}} \end{pmatrix}$

$y = aX \oplus W$  with  $\mathbb{Z}$ -action  
 $V \oplus bX$  in order to complete you must give an isom  $W \cong V$

722

Choose basis  $\xi_0$  forIs  $Y$  a fund. domain?Basically you have  $Y = aX \oplus W$   
 $= V \oplus bX$ and you take ~~a copy~~ a copy of  $Y$   
at each vertex: ~~a copy~~  $\oplus z^n Y$ then you define  $u$  to be the isom ~~b\*~~:  $aX \rightarrow bX$   
direct sum with  $z$  times the isom of  $W$  with  $V$   
This means we have

$$\begin{matrix} X & V \\ \oplus = \oplus & \nearrow X \\ W & X \end{matrix} \quad \begin{matrix} \oplus \\ zW \end{matrix}$$

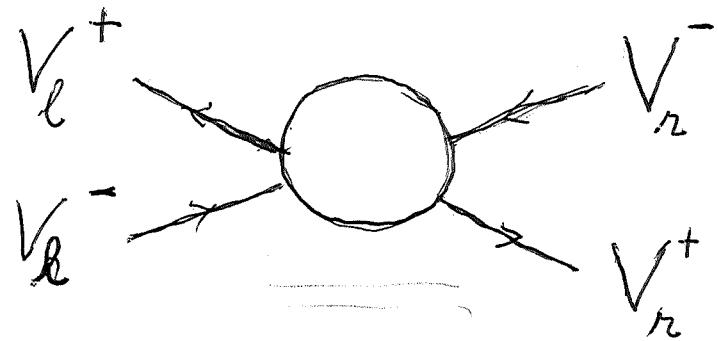
couple periodically.

Start carefully. The problem above is  
that it's one-dim.

Start again. You need a 2 port to begin and you propose one with  $X=0$ . So all you have is a  $U(1,1)$  matrix, between the left states and the right states. But then when you couple you need to identify the left and the right states. So in the end you have ~~constructed~~ constructed a flat bundle over the circle - cylinder construction. What's strange is that <sup>seem to</sup> we have a Hilbert space with unitary operator made out of ~~these~~ the sections of this bundle

723 Start with a unitary matrix having invertible transmission. This should be ~~Hermitian~~ the same space Hermitian  $\mathbb{H}_2$  graded a self isomorphism of the same space. Really you are looking at  $g \in \underline{\mathcal{U}(2)}$ , because you must identify the in and out spaces in order to proceed further. Presumably it should not matter if I conjugate by a diagonal matrix. ~~so that everything should depend ultimately~~

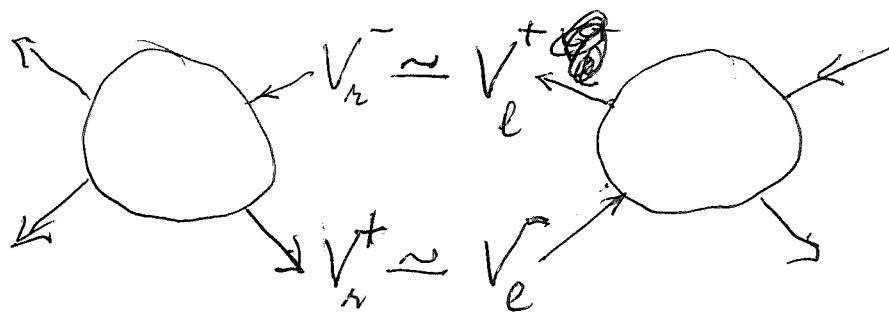
Start again with a 2-port indep of freq.  
left state space                      right state space



$$\left\| \begin{pmatrix} \xi^+ \\ \zeta \end{pmatrix} \right\|^2 + \left\| \begin{pmatrix} \xi^+ \\ \zeta \end{pmatrix} \right\|^2 = \left\| \begin{pmatrix} \xi^- \\ \zeta \end{pmatrix} \right\|^2 + \left\| \begin{pmatrix} \bar{\xi}^- \\ \bar{\zeta} \end{pmatrix} \right\|^2 \quad \text{cons. of eng}$$

$$-\|\xi_e^+\|^2 + \|\xi_e^-\|^2 = \|\xi_n^+\|^2 - \|\xi_n^-\|^2$$

Now couple these to form a chain

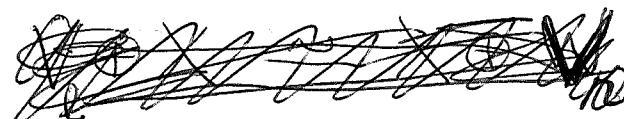


724 Now what you need is to construct the Hilbert space + unitary operator arising from this stuff.

Coupling two circuits yields another 2 port given by the product of transfer matrices.  $V^+, V^-$ . ~~Then~~

~~But~~ I've run into a problem, namely, ~~that~~ you seem to need an  $X$  around before iterating a 2 port yields an interesting Hilbert space and unitary operator

Consider the process of connecting 2 ports. Recall that a 2-port consists of a partial unitary



$$V^- \oplus X \simeq X \oplus V^+$$

together with decompositions  $V^+ = V_e^+ \oplus V_r^+$   $V^- = V_e^- \oplus V_r^-$

To couple these you take direct sum

$${}^1V_e^- \oplus {}^1X \oplus {}^1V_r^+ \simeq {}^1V_e^+ \oplus {}^1X \oplus {}^1V_r^+$$

$${}^2V_e^- \oplus {}^2X \oplus {}^2V_r^- \simeq {}^2V_e^+ \oplus {}^2X \oplus {}^2V_r^+$$

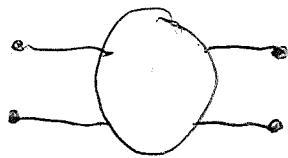
725 You need isomorphisms

$${}^1V_r^+ \simeq {}^2V_l^-$$

$${}^1V_r^- \simeq {}^2V_l^+$$

From this you should be able to get an isom.

$${}^1V_l^- \oplus {}^1X^0 \oplus {}^2X^0 \oplus {}^2V_r^+ \quad ??$$



What do you think happens? A 2-port is like a R.S. with 2 boundary components. A boundary component can be identified with a ~~pseudo~~ pseudo hemispherical plane.

Coupling requires ~~gluing~~ glueing boundary components corresponds to coupling transform matrices.

Somehow you have to understand what's going on. You need to do things ~~globally~~ on a line. The aim of ~~this~~ gluing is to obtain an analytic function of  $z$  whose zeroes are on the unit circle. This means essentially

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726 You want to construct a unitary op by glueing partial unitaries together. Finite dimensional partial unitary have degree =  $\dim X$ .

I need good examples. If I want inf. dim unitary then I need "infinite degree". So I need at least to start with  $\dim(X) = 1$  and glue copies of this together. Understand this carefully.

~~So I want to look at the simplest~~

So I want to self-couple a two port so as to obtain a ~~constant~~ unitary with  $\mathbb{Z}_2$ -symmetry. ~~symmetries~~ So start with the 2-port

$$Y = aX \oplus W = V \oplus bX$$

Here  $W, V$  have dim 2,  $X$  has dim 1, so  $V \oplus W$  is 1-diml. ~~But~~ But  $W, V$  are split.

~~Really~~  $V = V^+$ ,  $W = V^-$  where

$$V^+ = V_l^+ \oplus V_r^+ \quad V^- = V_l^- \oplus V_r^- \quad \text{and to } \text{couple}$$

couple one ~~gate~~ needs to identify the left and right ~~gates~~ gates:  $V_r^- \simeq V_l^+$   $V_r^+ \simeq V_l^-$

Now this situation may seem special, but you should be able to see a Hilbert space and unitary operator in all of this. ~~It's~~

Look: The Hilbert.

727 You start with  $X \rightarrow Y$  and want to couple two copies.

Consider Schrödinger equation periodic potential.

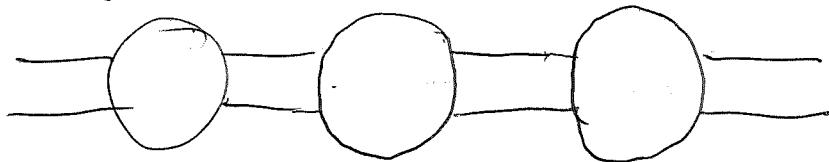
You need to understand the s.a. operator versus the transfer matrix, Bloch eigenfunctions, Q-Hall eff.

Let's see what you understand, or can construct.

You consider zitch The real

Try the following: Begin with a 2-port with invertible transmission, ~~what 2-port?~~

Do slowly. You take a 2-port



and couple it to itself periodically. Does this give you a Hilbert space & unitary operator commuting with  $\mathbb{Z}$  translation? I have been assuming this is obvious. There should not be a problem. Each port gives a partial unitary

$$V \oplus X \simeq X \oplus V^*$$

You then have a partial unitary:

$$\bigoplus \mathbb{Z}^n V^- \oplus \bigoplus \mathbb{Z}^n X \simeq \bigoplus \mathbb{Z}^n X \oplus \bigoplus \mathbb{Z}^n V^+$$

and you have an identification coming from the splitting  $V^- = V^- \oplus V^+$   $V^+ = V$

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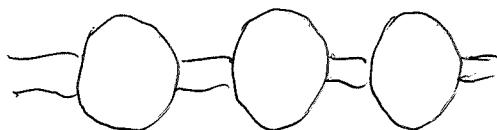
This works even if  $X = 0$ .

Somewhere here there is a way to handle ~~a~~ a constant unitary. You start with  $V^- \simeq V^+$

Feb 27. Coupling. ~~all those~~

You have to set up a proper notation.

The idea is to take a 2 port and connect it to itself ~~infinitely~~ periodically.



so as to obtain something with  $\mathbb{Z}$ -symmetry. This should be an analog of a Schrödinger equation with periodic potential.

Begin by coupling two copies:

$$V_l^- \oplus V_r^- \otimes X \xrightarrow{\quad} X \oplus V_l^+ \oplus V_r^+$$

$$zV_l^- \oplus zV_r^- \oplus zX \xrightarrow{2021} zX \oplus (zV_l^+ \oplus zV_r^+)$$

We are given  $V_r^+ \simeq zV_l^-$ ,  $V_r^- \simeq zV_l^+$   
Take direct sum you get

$$V_l^- \oplus zV_r^- \oplus L \oplus M \oplus X \oplus zX \simeq \\ X \oplus zX \oplus M \oplus L \oplus V_l^+ \oplus zV_r^+$$

729 Point: You ~~can't~~ take the direct sum of two <sup>partial</sup> unitaries, then extend the partial unitary by using the boundary identifications. The important thing to note is Notice that the ~~partial unitary~~ will be defined on  $X \oplus zX$   $\oplus$  the ~~gluing connecting the~~ connecting gate, ~~which is~~ which is  $V_l^+ \simeq zV_r^-$ ,  $\oplus V_r^- \simeq zV_l^+$ . Note that partial unitary is ~~too~~

so what do I do in general. You are trying to put something ~~over the disk~~ over a circle. You want shift operator

How to set things up over  $\mathbb{R} = \tilde{S}^1$ ?

Basically you put something over  $[0, 1]$  and give clutching data at the ends. What does this mean in our case. You start with the partial unitary

$$Y = aX \oplus V^+ = V^- \oplus bX \quad \begin{matrix} \text{(notice order} \\ \text{different from preceding} \\ \text{page)} \end{matrix}$$

$$V^+ = V_l^+ \oplus V_r^+ \quad V^- = V_l^- \oplus V_r^-$$

gluing:  $V_r^+ \simeq zV_l^-$ ,  $V_r^- \simeq zV_l^+$

I think you want

$$\begin{aligned} \oplus z^n Y &= (\oplus z^n aX) \oplus (\oplus z^n V^+) \\ &= (\oplus z^n V^-) \oplus (\oplus z^n bX) \end{aligned}$$

730 Does it have any eigenfunctions

The important point is that we produce a unitary operator because then it ~~should~~ might be possible to get a functional equation.

$$\begin{aligned} L^2(S^1, Y) &= L^2(S^1, aX) \oplus L^2(S^1, V_e^+) \oplus L^2(S^1, V_r^+) \\ &\quad \xrightarrow{\text{Zo (given iso)}} \quad \xleftarrow{\text{Zo (given iso)}} \end{aligned}$$
$$= L^2(S^1, V_e^-) \oplus L^2(S^1, V_r^-) \oplus L^2(S^1, bX)$$

take  $X=0$ . Then we have two orthogonal splittings

$$L^2(S^1, Y) = L^2(S^1, V_e^+) \oplus L^2(S^1, V_r^+)$$

$$= L^2(S^1, V_r^-) \oplus L^2(S^1, V_e^-)$$

$$Y = V_e^+ \oplus V_r^+ = V_r^- \oplus V_e^-$$

These are two orthogonal splittings and in addition you give isos.  $V_e^+ \xrightarrow{\sim} V_r^-$   $V_r^+ \xrightarrow{\sim} V_e^-$ . Current dimensions  $S^2 \times S^2 \times S^1 \times S^1$  dim=6. Act on it by  $U(2)$ . NO! Choose orth. bases for  $V_e^+$ ,  $V_r^+$ ; then

If you fix ~~the~~ the <sup>orth</sup> decomp.  $Y = V_e^+ \oplus V_r^+$ , then to give another orth. ~~as~~ decmp  $Y = V_e^- \oplus V_r^-$  together w <sup>unit</sup> isos.  $V_e^+ \xrightarrow{\sim} V_r^-$   $V_r^+ \xrightarrow{\sim} V_e^-$  is the same as

731 ~~Weyl's rule~~ so take a  $2 \times 2$  unitary matrix  $\begin{pmatrix} a & b \\ c & d \end{pmatrix}$  and then we get ~~a~~ a unitary in  $L^2(\mathbb{S})^{\oplus 2}$  commuting with  $z$ , should be a <sup>unitary</sup> matrix function, namely  $\begin{pmatrix} az & b \\ c & dz^{-1} \end{pmatrix}$ . Since this commutes with  $z$  there should be simultaneous ~~different~~ "eigenvectors" essentially, namely eigenvectors for the above matrix.

$$\lambda^2 - (az + dz^{-1})\lambda + (ad - bc) = 0$$

~~Suppose~~ Suppose  $ad - bc = 1$ .  $\begin{pmatrix} a & b \\ -b & \bar{a} \end{pmatrix}$

$$\lambda^2 - (az + \bar{a}z^{-1})\lambda + 1 = 0.$$

~~2 cos(θ)~~

Suppose  $a = 1$ . ~~Then~~ The roots are  $\lambda = z, z^{-1}$ .

The unitary  $\begin{pmatrix} az & b \\ c & dz^{-1} \end{pmatrix} = \begin{pmatrix} z & 0 \\ 0 & z^{-1} \end{pmatrix}$

Next look at the <sup>1-part given by the</sup> half-line. What is the ~~response~~ response function? Use the same eigenfunctions.

Wait first take  $0 < a < 1$ . Then for ~~the~~

$|\lambda| \leq 1$   $\frac{az + \bar{a}z^{-1}}{2}$  is real in  $[-a, a]$

and  $\lambda = e^{\pm i\theta}$  where  $\cos \theta = \frac{az + \bar{a}z^{-1}}{2}$ .

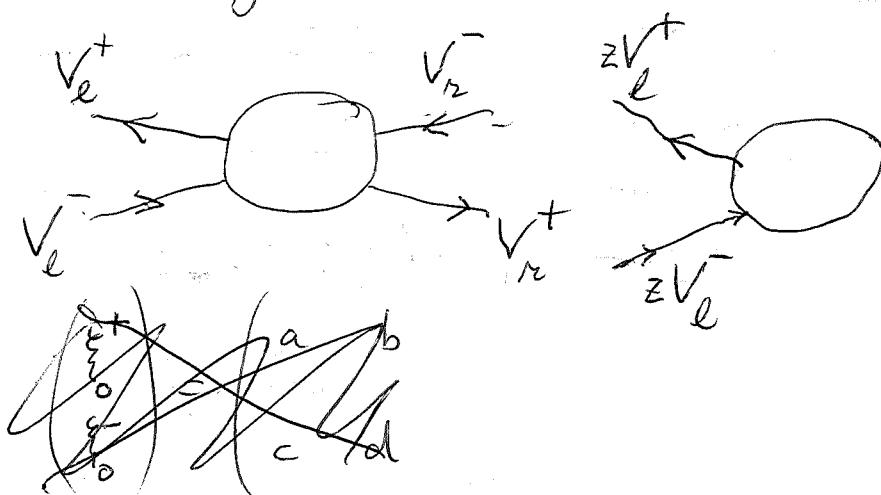
As  $z$  ranges over  $S^1$ ,

732 We have a curve:  $(\lambda, z)$  satisfying

$$\lambda^2 - (az + \bar{a}\bar{z}^{-1})\lambda + 1 = 0$$

describing the simultaneous "eigenfunctions" for  $z$  and  $\begin{pmatrix} az & b \\ -\bar{b} & \bar{a}\bar{z}^{-1} \end{pmatrix}$ .

~~Below~~ To finish the picture you should check results via the transfer matrix. You have the unitary matrix



$$\begin{pmatrix} \xi_e^+ \\ \xi_e^- \\ \xi_r^+ \\ \xi_r^- \end{pmatrix} = \begin{pmatrix} a & b \\ -\bar{b} & \bar{a} \end{pmatrix} \begin{pmatrix} \xi_r^+ \\ \xi_r^- \end{pmatrix}$$

$$\xi_e^+ = a \xi_r^+ + b \xi_r^- \quad \Rightarrow \quad \xi_e^+ = \left( a + \frac{|b|^2}{\bar{a}} \right) \xi_r^+ + \frac{b}{\bar{a}} \xi_r^-$$

$$\xi_r^+ = -\bar{b} \xi_r^+ + \bar{a} \xi_r^- \quad \Rightarrow \quad \xi_r^+ = \frac{1}{\bar{a}} \xi_r^+ + \frac{\bar{b}}{\bar{a}} \xi_r^-$$

$$\begin{pmatrix} \xi_e^- \\ \xi_r^+ \\ \xi_r^- \end{pmatrix} = \begin{pmatrix} \frac{1}{\bar{a}} & \frac{\bar{b}}{\bar{a}} \\ \frac{b}{\bar{a}} & \frac{1}{\bar{a}} \end{pmatrix} \begin{pmatrix} \xi_r^+ \\ \xi_r^- \end{pmatrix} = \begin{pmatrix} \frac{1}{\bar{a}} & \frac{\bar{b}}{\bar{a}} \\ \frac{b}{\bar{a}} & \frac{1}{\bar{a}} \end{pmatrix} \begin{pmatrix} z \xi_e^- \\ z \xi_r^- \end{pmatrix}$$

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 $\frac{1-|b|^2}{\bar{a}^2} = \frac{a}{\bar{a}^2} = \frac{a}{\bar{a}^2}$

$$\begin{pmatrix} \frac{z}{\bar{a}} & \frac{1}{\bar{a}} z^{-1} \\ \frac{b}{\bar{a}} z & \frac{1}{\bar{a}} z^{-1} \end{pmatrix}$$

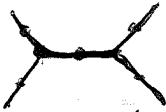
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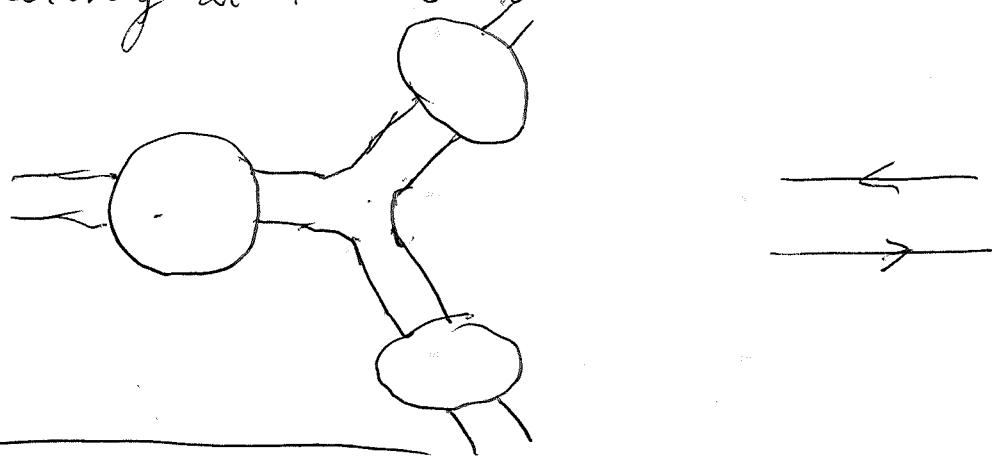
$$\lambda^2 - \left( \frac{z}{\bar{a}} + \frac{\bar{z}}{\bar{a}} \right) \lambda + \frac{a}{\bar{a}} = 0$$

733. I have to move on to Roe's stuff.

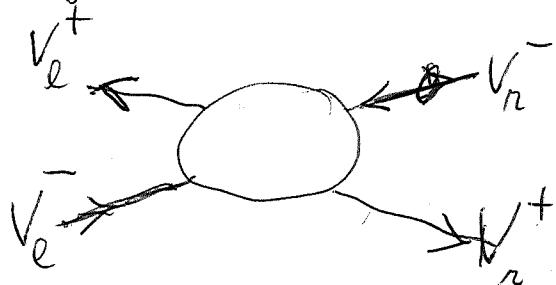
$\Gamma = PSL_2(\mathbb{Z})$  acting on tree



You want the simplest example with scattering  
Use cyclic ordering at the 3 vertices.



Go on. You need a 2 port with  $\mathbb{Z}/2$  acting



$$Y = V^- \oplus aX \xrightarrow{=} bX \oplus V^+$$

I think you want  $\mathbb{Z}/2$  to act on  $Y$  preserving  $X$ ,  $a$ , and  $b$ , hence also  $V^-$  and  $V^+$ . It should interchange  $V_r^+$  and  $V_l^+$ , also  $V_l^-$  and  $V_r^-$ . If  $X$  is zero you have simply  $Y = V^- = V^+$  hence ~~a unitary operator~~ two polarizations of  $Y$  namely  $V_l^- \oplus V_r^-$  and  $V_l^+ \oplus V_r^+$

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# 734 more 2-port stuff

$$Y = V^- = V^+$$

$$Y = V_e^- \oplus V_r^- = V_e^+ \oplus V_r^+$$

want  $\mathbb{Z}/2$  action, that  flips left and right. Can suppose  $Y = \mathbb{C}^2$  with  $\sigma = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$ . Ask when  $L \perp \sigma L$ ?

Suppose  $L = \mathbb{C}\begin{pmatrix} 1 \\ k \end{pmatrix}$ , then  $\sigma L = \mathbb{C}\begin{pmatrix} 1 \\ -k \end{pmatrix}$ , these are  $\perp$  where  $\begin{pmatrix} 1 \\ k \end{pmatrix}^* \begin{pmatrix} 1 \\ -k \end{pmatrix} = (1 \bar{k}) \begin{pmatrix} 1 \\ -k \end{pmatrix} = 1 - |k|^2 = 0$

so the possibilities are described by  $S^1$ .

It's time to examine ~~examine~~ 2-ports with  $\mathbb{Z}/2$  symmetric. Want  ~~$X \xrightarrow[a]{b} Y$~~  unitary reps of  $\mathbb{Z}/2$ . ~~These are split into two~~  
A 2-port is more than a partial unitary - It's like a manifold with two boundary components

First discuss  $X=0$ . Have  $Y = V^- = V^+$  where  $V^- = V_e^- \oplus V_r^-$ ,  $V^+ = V_e^+ \oplus V_r^+$ . Thus  $Y$  is a space with 2 <sup>orth</sup> decompositions. Have  $\mathbb{Z}/2$  action on  $Y$  preserving these decompositions

7.35 Suppose  $\mathcal{Y} = \mathbb{C}^2$  with  $\sigma = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$

Let  $L = \mathbb{C} \begin{pmatrix} x \\ y \end{pmatrix}$  be such that  $L$  and  $\sigma L$  are  $\perp$ .

$$\begin{pmatrix} x \\ y \end{pmatrix}^* \begin{pmatrix} y \\ x \end{pmatrix} = x^*y + y^*x = \bar{x}y + \bar{y}x = 0$$

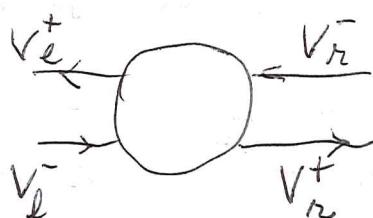
e.g.  $\begin{pmatrix} 1 \\ 0 \end{pmatrix} \begin{pmatrix} 0 \\ 1 \end{pmatrix}$ : Can suppose  $x$  real, if  $x \neq 0$   
then  $y$  is purely imag. So dealing with

$$L = \cancel{\begin{pmatrix} x \\ y \end{pmatrix}} \quad L = \begin{pmatrix} 1 \\ it \end{pmatrix} \subset \mathbb{C} \quad t \in \mathbb{R}$$

$$L, L^\perp = \begin{pmatrix} \cos \theta & \cancel{\sin \theta} \\ \cancel{\sin \theta} & \cos \theta \end{pmatrix}$$

unitary matrix  $\begin{pmatrix} a & b \\ -\bar{b} & \bar{a} \end{pmatrix}$  with  $a$  real  
and  $b \in i\mathbb{R}$ .

Life goes on.



$$\begin{pmatrix} \xi^+ \\ \xi^- \\ \xi^+ \\ \xi^- \end{pmatrix} = \begin{pmatrix} a & b \\ c & d \end{pmatrix} \begin{pmatrix} \xi^- \\ \xi^+ \end{pmatrix}$$

$$\begin{pmatrix} d & c \\ b & a \end{pmatrix} = \begin{pmatrix} a & b \\ c & d \end{pmatrix}$$

inv. under conj  
by  $\begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$

$$\begin{pmatrix} a & b \\ b & a \end{pmatrix}$$

and unitary means

$$|a|^2 + |b|^2 = 1$$

$$ab + b\bar{a} = 0.$$

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736 Let's now understand the effect of  $X$ . We are given

$$Y = V \ominus bX = \boxed{\phantom{00}}$$

(Still you have to get  $a, b$  straight.)

Let's review the scattering business.

$$Y = aX \oplus \text{Ker}(a^*) = \text{Ker}(b^*) \oplus bX$$

$$= \overline{(aX + bX)} \oplus (\text{Ker}(a^*) \cap \text{Ker}(b^*))$$

$$\begin{aligned} \|ax + bx'\|^2 &= \|x\|^2 + \|x'\|^2 + (ax, bx') + (bx', ax) \\ &= \|x\|^2 + \|x'\|^2 + (b^*ax, x') + (x', b^*a x) \end{aligned}$$

Put  $h = b^*a : X \rightarrow X$ . Then  $aX + bX$  can be recovered from  $X, h$  as the completion of  $X \oplus X$  w.t.  $\|x + x'\|^2 = \|x\|^2 + \|x'\|^2 + (hx, x') + (x, hx)$

~~$$\cancel{\|x\|^2 + \|x'\|^2 + (hx, x') + (x, hx)}$$~~

$$\|hx\|^2 - \|hx\|^2 = \|x\|^2 - \|hx\|^2$$

$$= \|x' + hx\|^2 + \|x\|^2 - \|hx\|^2$$

so find  $\overline{aX + bX} \simeq X \oplus \sqrt{1-h^*h} X$   
 $ax + bx' \mapsto (hx + x') \oplus \sqrt{1-h^*h} x$

Similarly find  $\overline{aX + bX} \simeq X \oplus \sqrt{1-hh^*} X$   
 $ax + bx' \mapsto (x + h^*x) \oplus \sqrt{1-hh^*} x$

Next discuss the scattering operator. Motivation  
~~eigenvalues~~ eigenvectors for partial unitary.

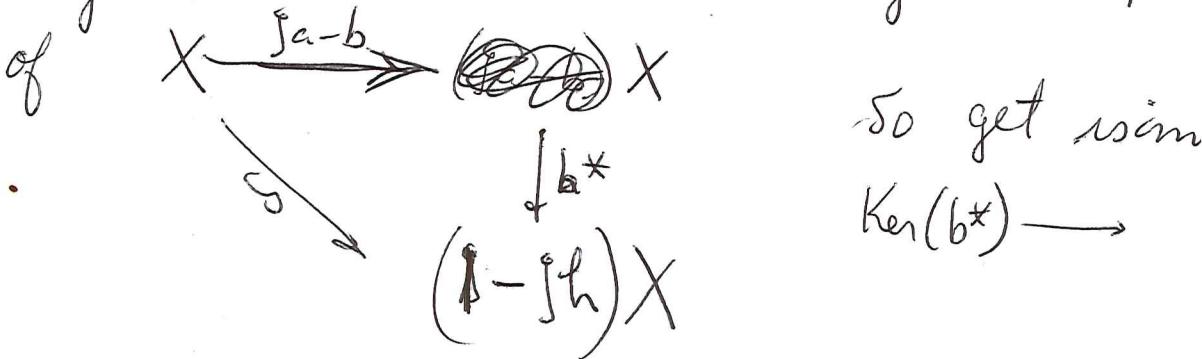
737. Scattering operator?  $h = b^*a$

~~$a - \bar{f}^{-1}b : X \rightarrow Y$~~

$$\begin{aligned}\|(\bar{a} - \bar{f}^{-1}b)x\|^2 &= (x, (\bar{a}^* - \bar{f}^{-1}b^*)(\bar{a} - \bar{f}^{-1}b)x) \\ &= (x, ((1 + |\bar{f}|^2) - \bar{f}^{-1}h - \bar{f}^{-1}h^*)x)\end{aligned}$$

$$\begin{aligned}\|\bar{f}a - b\|^2 &= \|\bar{f}ax\|^2 + \|x\|^2 - (\bar{f}ax, bx) - (bx, \bar{f}ax) \\ &= \|x\|^2 + |\bar{f}|^2 \|x\|^2 - (\bar{f}hx, x) - (x, \bar{f}hx) \\ &\quad + \|\bar{f}hx\|^2 - \|\bar{f}hx\|^2 \\ &= \|x - \bar{f}hx\|^2 + \underbrace{|\bar{f}|^2(\|x\|^2 - \|hx\|^2)}_{\geq 0} \\ \|\bar{f}x - \bar{f}hx\| &\geq \|x\| - \|\bar{f}hx\| \geq \|x\| - |\bar{f}|\|x\|\end{aligned}$$

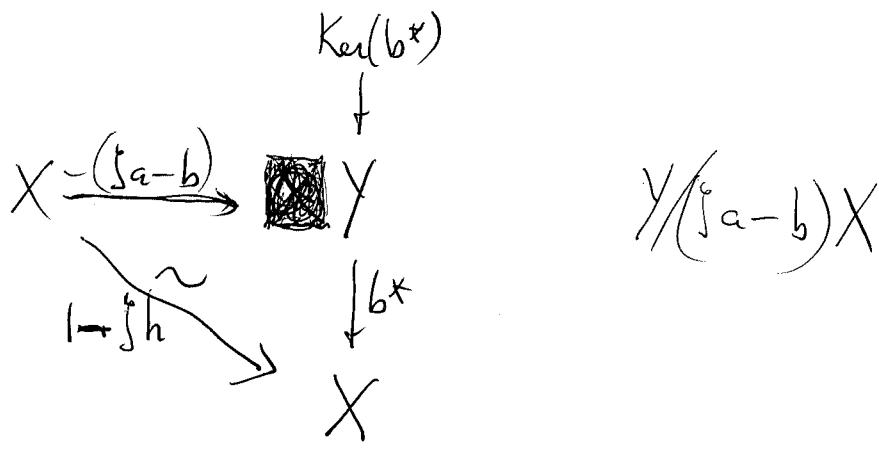
so if  $|\bar{f}| < 1$   $\bar{f}a - b$  is invertible no injective with closed image: top.  $\sim$



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Point.  $Y = (f_a - b)X \oplus \text{Ker}(b^*)$ .  $|S| < 1$ .

What are you doing? You have a partial unitary  $X \xrightarrow[a]{b} Y$ . It can be extended to a unitary iff  $\text{Ker}(a^*)$ ,  $\text{Ker}(b^*)$  are isom. You don't expect a scattering operator except in this case. How can you expect a unitary isom ~~between~~ between  $\text{Ker}(a^*)$ ,  $\text{Ker}(b^*)$  otherwise.

How do you actually construct the scattering operator. Simplest method is to couple  $Y$  to a transmission line.

Try to organize the  $\Gamma$  case. Straighten out the Hilbert space and unitary operator I want to ~~look~~ look at the case where both ports have zero  $X$ . Remember that when you connect you identify ~~two~~<sup>two</sup> gates, but ~~they became~~ ~~just~~ one copy appears in the Hilbert space