

In finite dimensions  $\exists$  formulas which enable one to evaluate integrals over fixed point submanifolds.

Finite-dim: symplectic manifold  $M^{2n}$  compact, action of  $S^1$   
(Duistermaat - Heckman)

~~Let~~  $\omega$  symplectic form,  $\frac{\omega^n}{n!} \neq 0$  volume  
 $X$  generator of  $S^1$  action.

$$dH = i(X)\omega$$

$H$  Hamiltonian

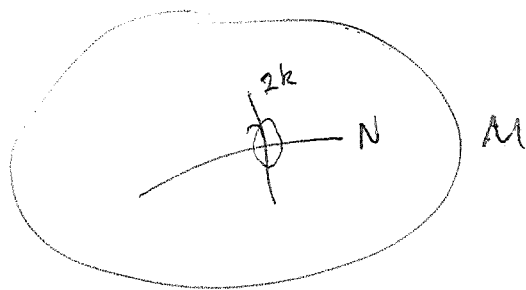
$$\int e^{-tH} \frac{\omega^n}{n!}$$

(Think of this formally as  $\int_M e^{\omega - tH}$ )

Thm:  $\int_M e^{-tH} \frac{\omega^n}{n!} = \sum_P \frac{e^{-tH(P)}}{t^n \prod_{j=1}^n m_j^P}$  isolated fixpts.

where  $z_1^{m_1}, \dots, z_n^{m_n}$  are the weights of  $S^1$  on  $T_P(M)$ .

Atiyah-Bott gave a cohomology proof.



$$c(M) = \prod_j (1 + \alpha_j)$$

where the  $\alpha_j$  are formal roots.

$$\int_M e^{-tH} \frac{\omega^n}{n!} = \sum_N \int_N \frac{e^{-tH(N)} e^\omega}{\prod_{j=1}^k (t m_j + \alpha_j)}$$

$\omega$  can be allowed to degenerate, except  $\omega^n = 0$   
~~is~~ an  $S \subset M$ , and defines  $[S] \in H^1(M, \mathbb{Z}_2)$ .  
 Thus this represents  $\omega_1$  of  $M$ .

Assume  $M$  orientable.

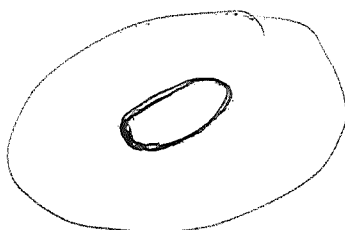
Given a Riemannian metric then  $\omega$  defines skew-adjoint  $A = -A^*$ .  $\det A = (\text{Pf} A)^n$ .

$$\frac{\omega^n}{n!} = (\text{Pf} A) \cdot \text{Riemannian volume}$$

$$\mathcal{QM} = \Omega M = \text{Map}(S^1, M)$$

$\dim M = 2n$  compact oriented with Riem. metric

$$\varphi: S^1 \hookrightarrow M$$



$$T_\varphi = \text{all vector fields along } S^1 = \Gamma(\varphi^* T_M)$$

has inner product  $\int_{S^1} \xi \cdot \eta$

$\exists$  natural skew-adjoint operator  $\nabla$  cov. deriv.

This form can be degenerate e.g. if geodesic.

$S^1$  acts on  $\Omega M$ . preserves  $\omega$ .

$$\text{Hamiltonian} = \text{Energy of loop} = \frac{1}{2} \int |\nabla \varphi|^2$$

When is  $\Omega M$  orientable? Answer  $\Leftrightarrow M$  is a spin manifold.

Explanation: Look where the form  $\omega$  become degenerate and calculate the element of  $H^1(\Omega M, \mathbb{Z}_2)$  and you get  $\omega_2 \in H^2(M, \mathbb{Z}_2)$ . Alternatively can look at the  $\parallel$  transport or monodromy along  $\varphi$ , call it  $T_\varphi$ , determined up to conjugacy ~~in~~ in  $SO(2n)$ .  $\alpha_1, \dots, \alpha_n$  angles of  $T_\varphi$

$$\begin{aligned} \text{Lemma: } |\det(\nabla \varphi)| &= \det(1 - T) = \prod (1 - e^{i\alpha_j})(1 - e^{-i\alpha_j}) \\ &= \prod_j [2 \sin(\alpha_j/2)]^2 \end{aligned}$$

$$\therefore \text{Pf}(\nabla_\varphi) = \sqrt{\det(1-T)}$$

so the manifold will be orientable when this can be chosen coherently.

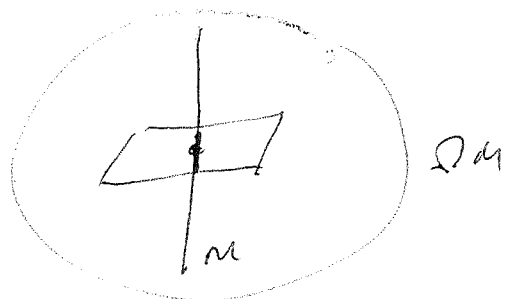
$$\begin{array}{ccc} \text{Spin}(2n) & \xrightarrow{\chi(s^+) - \chi(s^-)} & \text{is defined on} \\ \downarrow & \parallel & \text{Spin}(2n) \\ \text{SO}(2n) & \xrightarrow{\sqrt{\det(1-T)}} & \\ & \parallel & \\ & \xrightarrow{\det(1-T)} & \\ & \parallel & \\ & \xrightarrow{\sum (-1)^{\delta} \text{Tr}(\Lambda \delta T)} & \chi(\Lambda^{e\sigma}) - \chi(\Lambda^{o\sigma}) \end{array}$$

Thus  $M \text{ spin} \iff$  all monodromy can be lifted to spinor gp.  
 $\iff \text{Pf}(\nabla_\varphi)$  defined.

So everything is now set up to apply formula  
 Fixpts of  $S'$  on  $\Omega M$  are constant loops.

(a) Fixpt contribution

$$\int_M \frac{e^{-tH(x)} \mathbb{1} e^{w \cdot \mathbb{1}}}{\prod (t m_j + \alpha_j)}$$



normal directions are non-constant Fourier exp.  
 so denom. is

$$\prod_{n=1}^{\infty} \prod_j (tn + \alpha_j) = \frac{\text{fn on } S' \otimes TM}{\text{const}} = \frac{\sin \alpha_j/2}{\alpha_j/2}$$

$$\text{fixpt contribution} = \int_M \prod_j \frac{\alpha_j/2}{\sin \alpha_j/2} \hat{A}(M) \quad \text{where} \quad \text{Pont}(M) = \prod (1 - \alpha_j)^2$$

= index of Dirac operator

(b) Integration <sup>over</sup>  $\Omega M$ :

$$\int_{\Omega M} e^{-tH} \text{ Pfaff} \cdot \text{Riem volume}$$

$$= \int_{\Omega M} e^{-tH(\varphi)} \sqrt{1 - T_\varphi} d\varphi$$

$$= \int_{\Omega M} e^{-tH(\varphi)} \{ \text{Tr}(S_\varphi^+) - \text{Tr}(S_\varphi^-) \} d\varphi$$

Now remember  $\text{index}(D) = \text{Tr}(e^{-tD^*D}) - \text{Tr}(e^{-tDD^*})$   
and

$$\langle x | e^{-t\Delta} | y \rangle = \int_{\substack{\text{paths} \\ \Omega_t(x,y)}} e^{-H(\varphi)} d\varphi$$

Now when a vector bundle  $V$  is around with connection, then you expect  $\Delta_V = \nabla_V^* \nabla_V$

$$\langle x | e^{-t\Delta_V} | y \rangle \stackrel{\in \text{Hom}(V_y, V_x)}{=} \int e^{-H(\varphi)} T_\varphi d\varphi$$

fudge factors  $\nabla^* \nabla \neq D^* D$  by scalar curvature.  
so this <sub>scalar</sub> has to be put in.